# ZONAL POLYNOMIALS AND QUANTUM ANTISYMMETRIC MATRICES 

NAIHUAN JING ${ }^{1, a}$ AND ROBERT RAY ${ }^{1, b}$

${ }^{1}$ Department of Mathematics, North Carolina State University, Raleigh, NC 27695, USA and School of Sciences, South China University of Technology, Guangzhou, Guangdong 510640, China.
${ }^{a}$ E-mail: jing@math.ncsu.edu
${ }^{2}$ Department of Mathematics, Gonzaga University, Spokane, WA 99258, USA.
${ }^{b}$ E-mail: rayr@gonzaga.edu
$\|\|\|$


#### Abstract

We study the quantum symmetric spaces for quantum general linear groups modulo symplectic groups. We first determine the structure of the quotient quantum group and completely determine the quantum invariants. We then derive the characteristic property for quantum Phaffian as well as its role in the quantum invariant sub-ring. The spherical functions, viewed as Macdonald polynomials, are also studied as the quantum analog of zonal spherical polynomials.


## 1. Introduction

The regular representation of $G L(n, \mathbb{C})$ can be realized on the ring

$$
\begin{equation*}
A(X)=\mathbb{C}\left[x_{11}, x_{12}, \ldots, x_{n n}\right] \tag{1}
\end{equation*}
$$

where regular functions are polynomials of the matrix elements of the $n \times n$ matrices. It is well known that $A(X)$ is a completely reducible $G L(n, \mathbb{C})$ module and the associated irreducible polynomial sub-representations are parametrized by the set of partitions

$$
\begin{equation*}
P_{n}=\left\{\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right) \in \mathbb{Z}^{n} ; \lambda_{1} \geq \cdots \geq \lambda_{n} \geq 0\right\} . \tag{2}
\end{equation*}
$$

[^0]For a given $\lambda \in P_{n}$, there is a unique (up to isomorphism) irreducible representation $V(\lambda)$ with highest weight $\lambda$. Similarly, one considers the modules $A(\operatorname{Sym}(n))$ of the polynomials in the coordinates of the $n \times n$ symmetric matrix, and the module $A(\operatorname{Skew}(2 n))$ of the polynomials in the coordinates of the $2 n \times 2 n$ skew symmetric matrix. These representations decompose into the multiplicity free sums [4, 3]:

$$
\begin{align*}
A(\operatorname{Sym}(n)) & \simeq \bigoplus_{\lambda \in P_{n}} V(2 \lambda)  \tag{3}\\
A(\operatorname{Skew}(2 n)) & \simeq \bigoplus_{\lambda \in P_{2 n}} V\left(\lambda_{1}, \lambda_{1}, \lambda_{2}, \lambda_{2}, \ldots, \lambda_{n}, \lambda_{n}\right) \tag{4}
\end{align*}
$$

which are invariant under the action of $O(n, \mathbb{C})$ and $S p(2 n, \mathbb{C})$ respectively. As L. Hua first noticed and A. James later formulated that the $O_{n}$ and $S p_{2 n}$ invariants are one dimensional and the zonal spherical functions enjoy similar properties of Schur symmetric functions [4, 6, 13].

In the case of quantum analog of the symmetric pair of general linear groups and symplectic groups, Noumi and Letzter [15, 12] showed that the quantum spherical functions are indeed certain Macdonald symmetric functions by working on the quantum algebra of the enveloping algebras. We will study directly the quantum invariant ring as a subring of the quantum general linear group. As in [8], we compute the Hopf ideal of quantum invariants for the symplectic case using certain quadratic polynomials of matrix coefficients of quantum general linear groups.

A new feature in current work on quantum invariants is that we will study the important role played by Pfaffian as in the classical symplectic case. In the quantum case, the quantum Phaffian played an important role in the invariant theory as well [18]. We first give a closed form definition for the quantum Phaffian and study its representation-theoretic meaning in the quantum setting. Through this we are able to give an appropriate quantum analog of its relations with quantum determinant. As expected, quantum Phaffians enjoy similar properties as quantum determinant in the orthogonal case.

This paper is organized as follows. In Section 2 we first recall some basic facts of certain quantum algebras, in particular, we discuss a quantum deformation of $A(X)$ and the associated quantum version of $G L(n, \mathbb{C})$ as presented in Noumi, Yamada, and Mimachi [16] and we recall the quantized
universal enveloping algebra $U_{q}(\mathfrak{g l}(n, \mathbb{C}))$. In Section 3 we describe a quantum symplectic group, $S p_{q}(2 n, \mathbb{C})$. Since there does not seem to be a natural embedding of $S p_{q}(2 n, \mathbb{C})$ in $G L_{q}(2 n, \mathbb{C})$ we define $S p_{q}(2 n, \mathbb{C})$ invariants (left and right) in an infinitesimal manner, similar to an earlier construction by Jing and Yamada [8] of polynomial invariants for a quantum orthogonal group. These quantum symplectic invariants give us a quantum version of the regular functions of the antisymmetric matrices. In addition to defining the generators of these functions, we describe their relations and we discuss a construction of a quantum analog to the Pfaffian function.

We then describe a complete reduction of the $S p_{q}(2 n, \mathbb{C})$ invariant spaces (left and right) into irreducible modules and we follow with a construction and characterization of the associated bi-invariant space and its basis of zonal polynomials. In the last section, a connection between the zonal polynomials and certain Macdonald polynomials is discussed.

## 2. Quantum Groups

Quantum groups are defined as certain one-parameter deformations of the algebra of algebraic functions on simple Lie groups [17]. In other words, we will describe $A_{q}(X)$ to be like the classical algebra $A(X)$, except with noncommuting relations imposed upon its generators. Throughout the paper we will let $q$ be a complex number and for $q \neq 1$ we require that $q$ not be a root of unity.

## 2.1. $A_{q}(X), A(G)$ and $G L_{q}(n, \mathbb{C})$

We first define the algebra of functions $A_{q}(X)$ on $X=\operatorname{Mat}_{q}(n, \mathbb{C})$ as a noncommutative $\mathbb{C}$-algebra

$$
\begin{equation*}
A_{q}(X)=\mathbb{C}_{q}\left[x_{11}, x_{12}, \ldots, x_{n n}\right] \tag{5}
\end{equation*}
$$

generated by $x_{11}, x_{12}, \ldots, x_{n, n}$ and with relations

$$
\begin{aligned}
x_{i k} x_{j k} & =q x_{j k} x_{i k}, \quad x_{k i} x_{k j}=q x_{k j} x_{k i}, \\
x_{i l} x_{j k} & =x_{j k} x_{i l}, \\
x_{i k} x_{j l} & -x_{j l} x_{i k}=\left(q-q^{-1}\right) x_{i l} x_{j k},
\end{aligned}
$$

where $i<j$ and $k<l$. The relations can be visualized by the diagram (see Figure (1) with a "square" of generators.


Figure 1: $A_{q}(X)$ Relations, $x \rightarrow y$ implies $x y=q y x$.
$A_{q}(X)$ is a bialgebra using the same coproduct and counit maps as defined on $A(X)$, see [14].

Let $I$ and $J$ be two subsets of $\{1,2, \ldots, n\}$ with $\# I=\# J=r$ with ordered elements, i.e. $i_{1}<i_{2}<\ldots<i_{r} \in I$ and $j_{1}<j_{2}<\ldots j_{r} \in J$. The quantum $r$-minor determinants are defined as

$$
\begin{equation*}
\xi_{J}^{I}=\xi_{j_{1}, \ldots, j_{r}}^{i_{1}, \ldots, i_{r}}=\sum_{\sigma \in \mathfrak{S}_{r}}(-q)^{l(\sigma)} x_{i_{1} j_{\sigma(1)}} x_{i_{2} j_{\sigma(2)}} \ldots x_{i_{r} j_{\sigma(r)}} \tag{6}
\end{equation*}
$$

where $l(\sigma)$ denotes the number of pairs $(i, j)$ with $i<j$ and $\sigma(i)>\sigma(j)$. There is a unique quantum $n$-minor determinant, and it is denoted by $\operatorname{det}_{q}$ [8]. We define the algebra of regular functions $A(G)$ on the quantum group $G L_{q}(n, \mathbb{C})$ by adjoining $\operatorname{det}_{q}^{-1}$ to $A_{q}(X)$

$$
\begin{equation*}
A(G)=\left[x_{11}, x_{12}, x_{13}, \ldots, x_{n n}, \operatorname{det}_{q}^{-1}\right] \tag{7}
\end{equation*}
$$

Then, $G L_{q}(n, \mathbb{C})$ is defined as the spectrum of the Hopf algebra algebra $A(G)$, i.e.

$$
\begin{equation*}
G L_{q}(n, \mathbb{C})=\operatorname{Spec}(A(G)) \tag{8}
\end{equation*}
$$

one usually refers to $G L(n, \mathbb{C})$ simply as $A(G)$.
In addition to the relations of $A_{q}(X), A(G)$ also has the following relations 16]

$$
\begin{align*}
x_{i j} \cdot \operatorname{det}_{q}^{-1} & =\operatorname{det}_{q}^{-1} \cdot x_{i j}  \tag{9}\\
\operatorname{det}_{q}^{-1} \cdot \operatorname{det}_{q} & =\operatorname{det}_{q} \cdot \operatorname{det}_{q}^{-1}=1 \tag{10}
\end{align*}
$$

This allows us to define the algebra morphism $S: A(G) \rightarrow A(G)$ by

$$
\begin{equation*}
S\left(x_{i j}\right)=(-q)^{i-j} \xi_{\hat{i}}^{\hat{j}} \cdot \operatorname{det}_{q}^{-1} \quad 1 \leq i, j \leq n \tag{11}
\end{equation*}
$$

where $\hat{k}=\{1, \ldots, k-1, k+1, \ldots, n\} . S$ is the antipode for $A(G)$ and makes $A(G)$ a Hopf algebra.

### 2.2. Additional Quantum Groups

In addition to the above mentioned quantum groups, we need some additional subgroups of $G=G L_{q}(n, \mathbb{C})$.

The diagonal subgroup $H_{n}$ of $G L_{q}(n, \mathbb{C})$ is defined by its regular functions

$$
\begin{equation*}
A\left(H_{n}\right)=\mathbb{C}\left[t_{1}, t_{1}^{-1}, \ldots, t_{n}, t_{n}^{-1}\right] . \tag{12}
\end{equation*}
$$

Associated with this commutative Hopf algebra, we have the restriction map $\pi_{H}: A(G) \rightarrow A\left(H_{n}\right)$ defined by

$$
\begin{equation*}
\pi_{H}\left(x_{i j}\right)=\delta_{i, j} t_{i} \tag{13}
\end{equation*}
$$

The Borel subgroups $B_{+}$and $B_{-}$of $G L_{q}(n, \mathbb{C})$ consist of the upper and lower triangular matrices and are defined in terms of their associated Hopf algebras

$$
\begin{array}{ll}
A\left(B_{+}\right)=\mathbb{C}\left[b_{i j}\right], & i \leq j, \\
A\left(B_{-}\right)=\mathbb{C}\left[b_{i j}\right], & i \geq j . \tag{15}
\end{array}
$$

These algebras have relations induced from $A(G)$ and we note that the diagonal elements $b_{11}, \ldots, b_{n n}$ commute with each other, [16]. With each of these Hopf algebras we define the restrictions maps $\pi_{B_{+}}: A(G) \rightarrow A\left(B_{+}\right)$ and $\pi_{B_{-}}: A(G) \rightarrow A\left(B_{-}\right)$respectively by

$$
\begin{align*}
& \pi_{B_{+}}\left(x_{i j}\right)=\left\{\begin{array}{cc}
b_{i j}, & (1 \leq i \leq j \leq n) \\
0, & (i>j)
\end{array}\right.  \tag{16}\\
& \pi_{B_{-}}\left(x_{i j}\right)=\left\{\begin{array}{cc}
b_{i j}, & (1 \leq j \leq i \leq n) \\
0, & (j>i)
\end{array}\right. \tag{17}
\end{align*}
$$

### 2.3. Enveloping Algebra $U_{q}(\mathfrak{g})$

We recall the quantum universal enveloping algebra $U_{q}(\mathfrak{g})$ of $\mathfrak{g}=\mathfrak{g l}(n, \mathbb{C})$ or rather $\mathfrak{s l}(n, \mathbb{C})[9]$. Let $L_{n}$ be the free $\mathbb{Z}$-module of rank $n$ with the canonical basis $\left\{\epsilon_{1}, \ldots, \epsilon_{n}\right\}$, i.e. $L_{n}=\bigoplus_{k=1}^{n} \mathbb{Z} \epsilon_{k}$, endowed with the symmetric bilinear form $\left\langle\epsilon_{i}, \epsilon_{j}\right\rangle=\delta_{i j}$. We will define $\alpha_{k}=\epsilon_{k}-\epsilon_{k+1}$. Additionally, we will identify a partition $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right) \in P_{n}$ with $\lambda_{1} \epsilon_{1}+\cdots+\lambda_{n} \epsilon_{n} \in L_{n}$. We will refer to such an element of $L_{n}$ as a dominant integral weight. The fundamental weights are defined by $\Lambda_{k}=\epsilon_{1}+\cdots+\epsilon_{k}$ (see [8]). Now we define $U_{q}(\mathfrak{g})$ as the $\mathbb{C}$-algebra with generators $e_{k}, f_{k}(1 \leq k<n)$ and $q^{\lambda}\left(\lambda \in \frac{1}{2} L_{n}\right)$ with the following relations [16]:

$$
\begin{align*}
& q^{0}=1, \quad q^{\lambda} q^{\mu}=q^{\lambda+\mu}  \tag{18}\\
& q^{\lambda} e_{k} q^{-\lambda}=q^{\left\langle\lambda, \alpha_{k}\right\rangle} e_{k} \quad(1 \leq k<n)  \tag{19}\\
& q^{\lambda} f_{k} q^{-\lambda}=q^{-\left\langle\lambda, \alpha_{k}\right\rangle} f_{k} \quad(1 \leq k<n)  \tag{20}\\
& e_{i} f_{j}-f_{j} e_{i}=\delta_{i j} \frac{q^{\alpha_{i}}-q^{-\alpha_{i}}}{q-q^{-1}} \quad(1 \leq i, j<n)  \tag{21}\\
& e_{i}^{2} e_{j}-\left(q+q^{-1}\right) e_{i} e_{j} e_{i}+e_{j} e_{i}^{2}=0 \quad(|i-j|=1)  \tag{22}\\
& f_{i}^{2} f_{j}-\left(q+q^{-1}\right) f_{i} f_{j} f_{i}+f_{j} f_{i}^{2}=0 \quad(|i-j|=1)  \tag{23}\\
& e_{i} e_{j}=e_{j} e_{i}, \quad f_{i} f_{j}=f_{j} f_{i} \quad(|i-j|>1) \tag{24}
\end{align*}
$$

We define a coproduct, $\Delta_{U}$, and a counit, $\varepsilon_{U}$, on the generators by

$$
\begin{align*}
& \Delta_{U}\left(q^{\lambda}\right)=q^{\lambda} \otimes q^{\lambda}, \quad \varepsilon\left(q^{\lambda}\right)=1,  \tag{25}\\
& \Delta_{U}\left(e_{k}\right)=e_{k} \otimes q^{-\alpha_{k} / 2}+q^{\alpha_{k} / 2} \otimes e_{k}, \quad \varepsilon\left(e_{k}\right)=0  \tag{26}\\
& \Delta_{U}\left(e_{k}\right)=f_{k} \otimes q^{-\alpha_{k} / 2}+q^{\alpha_{k} / 2} \otimes f_{k}, \quad \varepsilon\left(f_{k}\right)=0, \tag{27}
\end{align*}
$$

making $U_{q}(\mathfrak{g})$ a bialgebra. Additionally, with the antipode $S_{U}$ defined by

$$
\begin{align*}
& S_{U}\left(q^{\lambda}\right)=q^{-\lambda},  \tag{28}\\
& S_{U}\left(e_{k}\right)=-q^{-1} e_{k},  \tag{29}\\
& S_{U}\left(f_{k}\right)=-q f_{k} . \tag{30}
\end{align*}
$$

$U_{q}(\mathfrak{g})$ becomes a Hopf algebra.

## 2.5. $A(G), U_{q}(\mathfrak{g})$ Duality

There exists a well-known dual pairing of Hopf algebras $U_{q}(\mathfrak{g})$ and $A(G)$

$$
\begin{equation*}
a(\varphi) \in \mathbb{C}, \quad a \in U_{q}(\mathfrak{g}), \varphi \in A(G) \tag{31}
\end{equation*}
$$

satisfying the following relations:

$$
\begin{align*}
q^{\lambda}\left(x_{i j}\right) & =\delta_{i, j} q^{\left(\lambda, \varepsilon_{i}\right\rangle}, \quad \lambda \in \frac{1}{2} L_{n}, \quad 1 \leq i, j \leq n  \tag{32}\\
e_{k}\left(x_{i j}\right) & =\delta_{i, k} \delta_{j, k+1}, \quad 1 \leq i, j \leq n  \tag{33}\\
f_{k}\left(x_{i j}\right) & =\delta_{i, k+1} \delta_{j, k}, \quad 1 \leq i, j \leq n  \tag{34}\\
q^{\lambda}\left(\operatorname{det}_{q}^{m}\right) & =q^{m\left\langle\lambda, \varepsilon_{1}, \ldots, \varepsilon_{n}\right\rangle} \quad m \in \mathbb{Z}  \tag{35}\\
e_{k}\left(\operatorname{det}_{q}^{m}\right) & =f_{k}\left(\operatorname{det}_{q}^{m}\right)=0 \quad m \in \mathbb{Z} \tag{36}
\end{align*}
$$

We extend these to the rest of $U_{q}(\mathfrak{g})$ and $A(G)$ by

$$
\begin{align*}
a(\varphi \psi) & =\Delta_{U}(a)(\varphi \otimes \psi)  \tag{37}\\
a(1) & =\varepsilon_{U}(a)  \tag{38}\\
(a b)(\varphi) & =(a \otimes b) \Delta(\varphi)  \tag{39}\\
1(\varphi) & =\varepsilon(\varphi) \quad\left(a, b \in U_{q}(\mathfrak{g}), \quad \varphi, \psi \in A(G)\right) \tag{40}
\end{align*}
$$

Additionally, we have

$$
\begin{equation*}
S_{U}(a) \cdot \psi=a \cdot S(\psi) \quad a \in U_{q}(\mathfrak{g}), \psi \in A(G) \tag{41}
\end{equation*}
$$

These relations realize a duality between the two Hopf algebras and allows us to regard the elements of $U_{q}(\mathfrak{g})$ as linear functionals on $A(G)$ (see 16]). This duality allows any right $A(G)$-comodule $V$ (resp. left $A(G)$-comodule $W$ ) with structure map $R_{G}: V \rightarrow V \otimes A_{q}(G)\left(\right.$ resp. $\left.L_{G}: W \rightarrow A(G) \otimes W\right)$ to become a left (resp. right) $U_{q}(\mathfrak{g})$-module with the following defined action

$$
\begin{gather*}
a . v=(i d \otimes a) R_{G}(v), \quad a \in U_{q}(\mathfrak{g}), v \in V,  \tag{42}\\
w \cdot a=(a \otimes i d) L_{G}(v), \quad a \in U_{q}(\mathfrak{g}), w \in W . \tag{43}
\end{gather*}
$$

More specifically, we already know $A_{q}(X)$ is a completely reducible two-sided $A(G)$-comodule using the comultiplication, $\Delta$, as the comodule structure map. As such, it becomes a completely reducible left and right $U_{q}(\mathfrak{g})$-module
[16, 8]. We can describe the left module action of the generators of $U_{q}(\mathfrak{g})$ on the generators of $A_{q}(X)$ by

$$
\begin{align*}
q^{\lambda} \cdot x_{i j} & =x_{i j} q^{\left\langle\lambda, \varepsilon_{j}\right\rangle}  \tag{44}\\
e_{k} \cdot x_{i j} & =x_{i, j-1} \delta_{j, k+1}  \tag{45}\\
f_{k} \cdot x_{i j} & =x_{i, j+1} \delta_{j, k} \tag{46}
\end{align*}
$$

and the right module action as

$$
\begin{align*}
x_{i j} \cdot q^{\lambda} & =x_{i j} q^{\left\langle\lambda, \varepsilon_{i}\right\rangle},  \tag{47}\\
x_{i j} \cdot e_{k} & =x_{i+1, j} \delta_{k, i},  \tag{48}\\
x_{i j} \cdot f_{k} & =x_{i-1, j} \delta_{k+1, i} . \tag{49}
\end{align*}
$$

### 2.6. Relative Invariants

For an element $\lambda=\sum_{k=1}^{n} \lambda_{k} \epsilon_{k} \in L_{n}$, let $z^{\lambda}=\prod_{k=1}^{n} z_{k k}^{\lambda_{k}} \in A\left(B_{ \pm}\right)$and $t^{\lambda}=\prod_{k=1}^{n} t_{k}^{\lambda_{k}} \in A(H)$, we define the spaces of relative invariants with respect to the subgroups $B_{ \pm}$by (see $[16,8]$ )

$$
\begin{align*}
& A\left(G / B_{+} ; z^{\lambda}\right)=\left\{\varphi \in A(G) ;\left(i d \otimes \pi_{B_{+}}\right) \Delta(\varphi)=\varphi \otimes z^{\lambda}\right\}  \tag{50}\\
& A\left(B_{-} \backslash G ; z^{\lambda}\right)=\left\{\varphi \in A(G) ;\left(\pi_{B_{-}} \otimes i d\right) \Delta(\varphi)=z^{\lambda} \otimes \varphi\right\} \tag{51}
\end{align*}
$$

where the restrictions maps $\pi_{ \pm}: A(G) \rightarrow A_{q}\left(B_{ \pm}\right)$are defined by $\pi_{B_{+}}\left(x_{i j}\right)=$ $z_{i, j}(1 \leq i \leq j \leq n), \pi_{B_{+}}\left(x_{i j}\right)=0(i>j)$, and $\pi_{B_{-}}\left(x_{i j}\right)=z_{i, j}(1 \leq j \leq i \leq$ $n), \pi_{B_{-}}\left(x_{i j}\right)=0(i<j)$.
$A\left(G / B_{+} ; z^{\lambda}\right)\left(\right.$ resp. $\left.A\left(B_{-} \backslash G ; z^{\lambda}\right)\right)$ is a left (resp. right) $A(G)$-subcomodule of $A(G)$ with structure mapping $\Delta$. It is proved in [16] that, for a dominant integral weight $\lambda \in P_{n}$, the space $A\left(G / B_{+} ; z^{\lambda}\right)\left(\right.$ resp. $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ ) gives a realization of the irreducible left (resp. right) $A(G)$-subcomodule $V_{q}^{L}(\lambda)\left(\right.$ resp. $\left.V_{q}^{R}(\lambda)\right)$ of $A_{q}(X)$, with highest weight $\lambda$.

## 3. Spaces of $q$-Symplectic Invariants

## 3.1. $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$

Here we describe a subalgebra of $U_{q}(\mathfrak{g})$ that is a quantum deformation of $U(\mathfrak{s p}(2 n, \mathbb{C}))$. Relative to the standard $n$ dimensional representation of $U_{q}(\mathfrak{g})$, we identify the generators $e_{k}$ of $U_{q}(\mathfrak{g})$ with $E_{k, k+1}$ and $f_{k}$ with $E_{k+1, k}$. If we let $\lambda=\lambda_{1} \epsilon_{1}+\cdots+\lambda_{2 n} \epsilon_{2 n} \in \frac{1}{2} L_{2 n}$, then $q^{\lambda}$ is represented by

$$
\begin{equation*}
q^{\lambda_{1}} E_{11}+q^{\lambda_{2}} E_{22}+\cdots+q^{\lambda_{2} n} E_{2 n, 2 n} \tag{52}
\end{equation*}
$$

We may then inductively generate the other elements, $E_{i, j}$, where $|i-j|>1$, by

$$
\begin{equation*}
E_{i, j}=E_{i, k} E_{k, j}-E_{k, j} E_{i, k} \tag{53}
\end{equation*}
$$

where $i<k<j$ or $j<k<i$ and $E_{i, j}$ and $E_{j, i}$ are independent of our choice of $k$, see [8].

We define the subalgebra $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$ of $U_{q}(\mathfrak{g})$ as the subalgebra generated by the following elements:

$$
\begin{array}{rlrl}
s p_{e}(i, j) & =E_{2 i-1,2 j}+q^{2(i-j)} E_{2 j-1,2 i} & 1 \leq i \neq j \leq n \\
s p_{e}(i, i) & =E_{2 i-1,2 i} & 1 \leq i \leq n \\
s p_{f}(i, j) & =E_{2 i, 2 j-1}+q^{2(i-j)} E_{2 j, 2 i-1} & 1 \leq i \neq j \leq n \\
s p_{f}(i, i) & =E_{2 i, 2 i-1} & 1 \leq i \leq n \\
s p_{h}(i, j) & =E_{2 i-1,2 j-1}-q^{2(i-j)} E_{2 j, 2 i} & 1 \leq i, j \leq n
\end{array}
$$

with $i, j \leq n$. It can be directly shown that the elements of the form

$$
\begin{array}{r}
s p_{e}(j, j), s p_{f}(j, j), \quad \text { where } 1 \leq j \leq n, \\
s p_{e}(i, i+1), s p_{f}(i, i+1), \quad 1 \leq i \leq n-1 \tag{60}
\end{array}
$$

generate $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$.

## 3.2. $q$-Symplectic Invariants

For a given left (resp. right) $U_{q}(\mathfrak{g})$-module $V$ (resp. $W$ ) we define the
$q$-symplectic invariants by

$$
\begin{align*}
V^{K} & =\left\{v \in V ; s p_{e}(i, j) \cdot v=0, s p_{f}(i, j) \cdot v=0 \quad 1 \leq i, j \leq n\right\}  \tag{61}\\
{ }^{K} W & =\left\{w \in W ; w \cdot s p_{e}(i, j)=0, w \cdot s p_{f}(i, j)=0 \quad 1 \leq i, j \leq n\right\} \tag{62}
\end{align*}
$$

Using the fact that $A_{q}(X)$ is a two-sided $U_{q}(\mathfrak{g})$-module (see 42, 43) we define the left and right quantum symplectic invariants in $A_{q}(X)$ as

$$
\begin{align*}
& A_{q}(X)^{K}=\left\{\varphi \in A_{q}(X) ; s p_{e}(i, j) \cdot \varphi=0, s p_{f}(i, j) \cdot \varphi=0\right.  \tag{63}\\
& { }^{K} A_{q}(X)=\left\{\varphi \in A_{q}(X) ; \varphi \cdot s p_{e}(i, j)=0, \varphi \cdot s p_{f}(i, j)=0\right.  \tag{64}\\
& 1 \leq i, j \leq n\}
\end{align*}
$$

The spaces $A_{q}(X)^{K}$ and ${ }^{K} A_{q}(X)$ are subalgebras of $A_{q}(X)$. Additionally, we see that $A_{q}(X)^{K}$ is a left $A(G)$-subcomodule of $A_{q}(X)$ (similarly ${ }^{K} A_{q}(X)$ is a right $A(G)$-subcomodule of $\left.A_{q}(X)\right)$. Equivalently, $A_{q}(X)^{K}$ is a right $U_{q}(\mathfrak{g})$-submodule of $A_{q}(X)$ and ${ }^{K} A_{q}(X)$ is a left $U_{q}(\mathfrak{g})$-submodule of $A_{q}(X)$.

Definition 3.1. For $n \in \mathbb{Z}_{+}$even, the following quadratic elements of $A_{q}(X)$ may be defined

$$
\begin{align*}
z_{i, j}^{L} & =\sum_{k=1}^{n} q^{(i+j+1-4 k) / 2}\left(x_{i, 2 k-1} x_{j, 2 k}-q x_{i, 2 k} x_{j, 2 k-1}\right) \\
& =\sum_{k=1}^{n / 2} q^{(i+j+1-4 k) / 2} \xi_{2 k-1,2 k}^{i, j},  \tag{65}\\
z_{i, j}^{R} & =\sum_{k=1}^{n} q^{-(i+j+1-4 k) / 2}\left(x_{2 k-1, i} x_{2 k, j}-q x_{2 k, i} x_{2 k-1, j}\right) \\
& =\sum_{k=1}^{n / 2} q^{-(i+j+1-4 k) / 2} \xi_{i, j}^{2 k-1,2 k} . \tag{66}
\end{align*}
$$

Using the fact

$$
\begin{align*}
e_{k} \cdot \xi_{r, s}^{i, j} & =\delta_{k, r-1} \xi_{r-1, s}^{i, j}+\delta_{k, s-1} \xi_{r, s-1}^{i, j}  \tag{67}\\
f_{k} \cdot \xi_{r, s}^{i, j} & =\delta_{k, r} \xi_{r+1, s}^{i, j}+\delta_{k, s} \xi_{r, s+1}^{i, j} \tag{68}
\end{align*}
$$

it can be shown that $z_{i, j}^{L}$ (resp. $z_{i, j}^{R}$ ) are annihilated by $s p_{e}(k, k), s p_{e}(k, k+1)$, $s p_{f}(k, k)$ and $s p_{f}(k, k+1)$, which is sufficient to show they are annihilated by all $s p_{e}(k, l)$ and $s p_{f}(k, l)$ and therefore $z_{i, j}^{L} \in A_{q}(X)^{K}\left(\right.$ resp. $\left.z_{i, j}^{R} \in{ }^{K} A_{q}(X)\right)$

We denote the subalgebra of $A_{q}(X)^{K}$ (resp. $\left.{ }^{K} A_{q}(X)\right)$ by $A_{q}^{L}(\mathcal{A})$ (resp. $\left.A_{q}^{R}(\mathcal{A})\right)$ generated by $z_{i, j}^{L}$ (resp. $z_{i, j}^{R}$ ). $A_{q}^{L}(\mathcal{A})$ is a left $A(G)$-subcomodule of $A_{q}(X)^{K}$ and $A_{q}^{R}(\mathcal{A})$ is a right $A(G)$-subcomodule of $A_{q}(X)^{K}$.

Theorem 3.2. The algebras $A_{q}^{L}(\mathcal{A})$ and $A_{q}^{R}(\mathcal{A})$ are isomorphic to the algebra $A_{q}(\mathcal{A})$ generated by $z_{i, j}(1 \leq i, 1 \leq j)$ with the following relations:

$$
\begin{align*}
& z_{i, j}=-q^{-1} z_{j, i}  \tag{69}\\
& z_{i, l} z_{j, k}=z_{j, k} z_{i, l}  \tag{70}\\
& z_{i, j} z_{i, k}=q z_{i, k} z_{i, j}  \tag{71}\\
& z_{i, k} z_{j, l}-z_{j, l} z_{i, k}=\left(q-q^{-1}\right) z_{i, l} z_{j, k},  \tag{72}\\
& z_{i, j} z_{k, l}-z_{k, l} z_{i, j}=\left(q-q^{-1}\right) z_{i, k} z_{j, l}-q\left(q-q^{-1}\right) z_{i, l} z_{j, k}, \tag{73}
\end{align*}
$$

where $i<j<k<l$.

Using Eq. (72) we may rewrite Eq. (73) as

$$
\begin{equation*}
z_{i, j} z_{k, l}-z_{k, l} z_{i, j}=q z_{j, l} z_{i, k}-q^{-1} z_{i, k} z_{j, l} \tag{74}
\end{equation*}
$$

The definitions of these generators also imply

$$
\begin{equation*}
z_{i, i}=0 \tag{75}
\end{equation*}
$$

### 3.3. Quantum Antisymmetric Matrices

If we denote by $\mathcal{A}$, the vector space of $n \times n$ antisymmetric matrices with basis

$$
\begin{equation*}
B_{\mathcal{A}}=\left\{E_{i, j}-E_{j, i} \mid 1<i<j \leq n\right\} \tag{76}
\end{equation*}
$$

then $\operatorname{dim}(\mathcal{A})=n(n-1) / 2$. We observe that $\operatorname{Hom}_{\text {Alg }}\left(A_{q}(\mathcal{A}), \mathbb{C}\right)$ is the set of $n \times n$ matrices with restrictions imposed by the relations Eq. (69) and Eq.
(75). If we denote $\operatorname{Hom}_{\text {Alg }}\left(A_{q}(\mathcal{A}), \mathbb{C}\right)$ by $\mathcal{A}_{q}$, and treat it as a vector space (in other words we are ignoring multiplication) we see its basis is

$$
\begin{equation*}
B_{\mathcal{A}_{q}}=\left\{E_{i, j}-q E_{j, i} \mid 1<i<j \leq n\right\} \tag{77}
\end{equation*}
$$

where $\operatorname{dim}\left(\mathcal{A}_{q}\right)=n(n-1) / 2$ and we have $\mathcal{A}_{q} \simeq \mathcal{A}$ as vector spaces. We may think of $\mathcal{A}_{q}$ as the quantum analog of the antisymmetric matrices.

### 3.4. Quantum Pfaffian

If $A=\left(a_{i, j}\right) \in \operatorname{Mat}(2 n, \mathbb{C})$ is an antisymmetric matrix, it can be written as

$$
A=\left[\begin{array}{cccc}
0 & a_{1,2} & \cdots & a_{1,2 n}  \tag{78}\\
-a_{1,2} & 0 & \cdots & a_{2,2 n} \\
\vdots & \vdots & \ddots & \vdots \\
-a_{1,2 n} & -a_{2,2 n} & \cdots & 0
\end{array}\right]
$$

and there exists a polynomial $f$ in $\mathbb{Z}\left[x_{i j}\right]$ such that $f^{2}(A)=\operatorname{det}(A),[5]$. This polynomial is called the Pfaffian, denoted Pf, and we write

$$
\begin{equation*}
P f^{2}(A)=\operatorname{det}(A) \tag{79}
\end{equation*}
$$

Moreover, if $B=\left(b_{i, j}\right) \in \operatorname{Mat}(2 n, \mathbb{C})$ and we define $A$ by

$$
a_{i, j}=\operatorname{det}\left[\begin{array}{ll}
b_{i, 1} & b_{i, 2}  \tag{80}\\
b_{j, 1} & b_{j, 2}
\end{array}\right]+\operatorname{det}\left[\begin{array}{ll}
b_{i, 3} & b_{i, 4} \\
b_{j, 3} & b_{j, 4}
\end{array}\right]+\cdots+\operatorname{det}\left[\begin{array}{ll}
b_{i, 2 n-1} & b_{i, 2 n} \\
b_{j, 2 n-1} & b_{j, 2 n}
\end{array}\right]
$$

then $A$ is antisymmetric and we have $\operatorname{Pf}(A)=\operatorname{det}(B),[5]$.
To construct an explicit formula for $P f$ we can define an index set $\Pi$, consisting of all ordered, 2-partitions of $2 n$. In other words,

$$
\begin{equation*}
\Pi=\left\{\left(i_{1}, j_{1}\right)\left(i_{2}, j_{2}\right) \ldots\left(i_{n}, j_{n}\right) ; i_{k}<j_{k} \text { and } i_{k}<i_{k+1}\right\} \tag{81}
\end{equation*}
$$

For example, if $2 n=4$ we have

$$
\begin{equation*}
\Pi=\{(1,2)(3,4),(1,3)(2,4),(1,4)(2,3)\} \tag{82}
\end{equation*}
$$

We can associate the elements of $\Pi$ with elements of the symmetric group $\mathfrak{S}_{2 n}$ in the following manner

$$
\pi \sim\left[\begin{array}{cccccc}
1 & 2 & 3 & 4 & \cdots & 2 n  \tag{83}\\
i_{1} & j_{1} & i_{2} & j_{2} & \cdots & j_{n}
\end{array}\right] \in \mathfrak{S}_{2 n}
$$

for $\pi=\left\{\left(i_{1}, j_{1}\right)\left(i_{2}, j_{2}\right) \ldots\left(i_{n}, j_{n}\right)\right\}$. This allows us to define $\operatorname{sgn}(\pi)$ and $l(\pi)$. If $A=\left(a_{i, j}\right)$ is an antisymmetric matrix we can then write

$$
\begin{equation*}
\operatorname{Pf}(A)=\sum_{\pi \in \Pi} \operatorname{sgn}(\pi) a_{\pi}=\sum_{\pi \in \Pi} \operatorname{sgn}(\pi) a_{i_{1} j_{1}} a_{i_{2} j_{2}} \cdots a_{i_{n}, j_{n}} \tag{84}
\end{equation*}
$$

Example 3.3. As an example, when $2 n=4$

$$
\begin{equation*}
\operatorname{Pf}(A)=a_{1,2} a_{3,4}-a_{1,3} a_{2,4}+a_{1,4} a_{2,3} \tag{85}
\end{equation*}
$$

Before we construct a quantum analog of the Pfaffian, we note that the quantum antisymmetric generators $z_{i, j}^{L}$ (resp. $z_{i, j}^{R}$ ), defined by Eq. (65) (resp. Eq. (66)), are in fact quantum analogs of Eq. (80). Additionally, we have already noted that $Z=\left(z_{i, j}^{L}\right)$ is a quantum antisymmetric matrix with the relation $z_{i, j}^{L}=-\frac{1}{q} z_{j, i}^{L}$ for $i<j$. We now use the same index set $\Pi$, to define the quantum Pfaffian as

$$
\begin{equation*}
P f_{q}(Z)=\sum_{\pi \in \Pi}(-q)^{l(\pi)} z_{\pi}^{L}=\sum_{\pi \in \Pi}(-q)^{l(\pi)} z_{i_{1} j_{1}}^{L} z_{i_{2} j_{2}}^{L} \cdots z_{i_{n}, j_{n}}^{L} . \tag{86}
\end{equation*}
$$

Remark 3.4. An inductive definition of quantum Phaffian was given in [18]. One can show that our definition matches with Strickland's.

Example 3.5. As an example, when $2 n=4$

$$
\begin{equation*}
P f_{q}(Z)=z_{1,2}^{L} z_{3,4}^{L}-q z_{1,3}^{L} z_{2,4}^{L}+q^{2} z_{1,4}^{L} z_{2,3}^{L} \tag{87}
\end{equation*}
$$

Theorem 3.6. For every positive even $2 n, P f_{q}(Z)=\operatorname{det}_{q}(X)$.
Proof. To show this equality, we will prove that $P f_{q}$ is simultaneously a highest and lowest weight vector for the right action of $U_{q}(\mathfrak{g})$. This will show $P f_{q}$ to be a scalar multiple of $\left(\operatorname{det}_{q}\right)^{c}$ for some $c \in \mathbb{Z}_{+}$.

To begin, we let $k$ be a positive integer such that $1 \leq k<2 n$. Since the right action of generators of $U_{q}(\mathfrak{g})$ on products of elements of $A_{q}(X)$ can be
described by [8],

$$
\begin{align*}
& \phi \psi \cdot e_{k}=(\phi \otimes \psi) \cdot\left(e_{k} \otimes q^{-a_{k} / 2}+q^{a_{k} / 2} \otimes e_{k}\right)  \tag{88}\\
& \phi \psi \cdot f_{k}=(\phi \otimes \psi) \cdot\left(f_{k} \otimes q^{-a_{k} / 2}+q^{a_{k} / 2} \otimes f_{k}\right) \tag{89}
\end{align*}
$$

we may expand this notation to describe the following right action of $e_{k}$ on the components of $P f_{q}$ as

$$
\begin{align*}
& z_{a_{1} b_{1}}^{L} z_{a_{2} b_{2}}^{L} \cdots z_{a_{n / 2} b_{n / 2}}^{L} \cdot e_{k} \\
& \quad=\quad z_{a_{1} b_{1}}^{L} \cdot e_{k} \otimes z_{a_{2} b_{2}}^{L} \cdot q^{-\alpha_{k} / 2} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot q^{-\alpha_{k} / 2} \\
& \quad+z_{a_{1} b_{1}}^{L} \cdot q^{\alpha_{k} / 2} \otimes z_{a_{2} b_{2}}^{L} \cdot e_{k} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot q^{-\alpha_{k} / 2} \\
& \quad \vdots \\
& \quad+\quad z_{a_{1} b_{1}}^{L} \cdot q^{\alpha_{k} / 2} \otimes z_{a_{2} b_{2}}^{L} \cdot q^{\alpha_{k} / 2} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot e_{k} \tag{90}
\end{align*}
$$

and

$$
\begin{align*}
& z_{a_{1} b_{1}}^{L} z_{a_{2} b_{2}}^{L} \cdots z_{a_{n / 2} b_{n / 2}}^{L} \cdot f_{k} \\
& \quad=z_{a_{1} b_{1}}^{L} \cdot f_{k} \otimes z_{a_{2} b_{2}}^{L} \cdot q^{-\alpha_{k} / 2} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot q^{-\alpha_{k} / 2} \\
& \quad+z_{a_{1} b_{1}}^{L} \cdot q^{\alpha_{k} / 2} \otimes z_{a_{2} b_{2}}^{L} \cdot f_{k} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot q^{-\alpha_{k} / 2} \\
& \quad \vdots  \tag{91}\\
& \quad+\quad z_{a_{1} b_{1}}^{L} \cdot q^{\alpha_{k} / 2} \otimes z_{a_{2} b_{2}}^{L} \cdot q^{\alpha_{k} / 2} \otimes \cdots \otimes z_{a_{n} b_{n}}^{L} \cdot f_{k}
\end{align*}
$$

Additionally, each of these $A_{q}^{L}(\mathcal{A})$ generators is a sum of quantum 2minor determinants (see Eq. (65)) in which the indices $i$ and $j$ of $z_{i j}^{L}$ define the rows for each of these quantum 2-minor determinants . As such, the right action of $e_{k}$ and $f_{k}$ on these generators can be described by the following,

$$
\begin{align*}
z_{i, j}^{L} \cdot e_{k} & =q^{-1 / 2}\left(\delta_{i, k} z_{k+1, j}^{L}+\delta_{j, k} z_{j, k+1}^{L}\right)  \tag{92}\\
z_{i, j}^{L} \cdot f_{k} & =q^{1 / 2}\left(\delta_{i, k+1} z_{k, j}^{L}+\delta_{j, k+1} z_{j, k}^{L}\right) \tag{93}
\end{align*}
$$

and the right action of $q^{\alpha / 2}$ and $q^{-\alpha / 2}$ are described by

$$
\begin{align*}
z_{i, j}^{L} \cdot q^{\alpha_{k} / 2} & =q^{1 / 2\left(\delta_{i, k}-\delta_{i, k+1}+\delta_{j, k}-\delta_{j, k+1}\right)} z_{i, j}^{L}  \tag{94}\\
z_{i, j}^{L} \cdot q^{-\alpha_{k} / 2} & =q^{1 / 2\left(-\delta_{i, k}+\delta_{i, k+1}-\delta_{j, k}+\delta_{j, k+1}\right)} z_{i, j}^{L} \tag{95}
\end{align*}
$$

For example

$$
\begin{equation*}
z_{3,4}^{L} \cdot q^{\alpha_{4} / 2}=q^{1 / 2} z_{3,4}^{L} \tag{96}
\end{equation*}
$$

Before we give a detailed description of the action of $e_{k}$ on $P f_{q}$, we show how the components of $\Pi$ may be paired, relative to the value of $k$. Since the components of $P f_{q}$ are indexed by all ordered 2-partitions, this will allow us to group the components of $P f_{q}$ in a way that the right action of $e_{k}$ (and $f_{k}$ ) will annihilate the pairs.

We first fix $k \in \mathbb{Z}$ such that $1 \leq k<2 n$. Now if we choose any of the ordered 2-partitions, say $\pi=\left(a_{1}, b_{1}\right)\left(a_{2}, b_{2}\right) \cdots\left(a_{n}, b_{n}\right)$, it must have an index $r$, containing $k$ and an index $s$ containing $k+1$. In other words, there exist $r$ and $s$ such that

$$
\begin{equation*}
k \in\left(a_{r}, b_{r}\right) \quad \text { and } \quad k+1 \in\left(a_{s}, b_{s}\right) \tag{97}
\end{equation*}
$$

This fixes $r$ and $s$. Also contained in the $\left(a_{r}, b_{r}\right)$ and $\left(a_{s}, b_{s}\right)$ pairs are two other integers, $u$ and $v$ such that $u<v$. If it happens that $r=s$, in other words, there exists $\left(a_{r}, b_{r}\right)$ such that $\left(a_{r}, b_{r}\right)=(k, k+1)$ then we will not pair it with another 2-partition. We will show later how the right action of $e_{k}$ and $f_{k}$ already annihilate it.

Example 3.7. Suppose $2 n=8$ and we fix $k=5$. One of the ordered 2 partitions of $\Pi$ is $(1,3)(2,6)(4,8)(5,7)$. In this case we see that $r=4$ and $s=2$. We then designate $u=2$ and $v=7$.

Now, with $r$ and $s$ still fixed, and for the designated $u$ and $v$, there are precisely three possibilities describing how $k, k+1, u$ and $v$ can be ordered. These are:

$$
\begin{align*}
& k<k+1<u<v  \tag{98}\\
& u<k<k+1<v  \tag{99}\\
& u<v<k<k+1 \tag{100}
\end{align*}
$$

For each of these possibilities we have the following,

- $k<k+1<u<v$

In this case, if $r \neq s$, there is another 2-partition, $\hat{\pi}$ identical to $\pi$ except in the $r^{t h}$ and $s^{t h}$ pairs, $u$ and $v$ are switched.

$$
\begin{align*}
& \pi=\left(a_{1}, b_{1}\right) \cdots(k, u)(k+1, v) \cdots\left(a_{n}, b_{n}\right)  \tag{101}\\
& \hat{\pi}=\left(a_{1}, b_{1}\right) \cdots(k, v)(k+1, u) \cdots\left(a_{n}, b_{n}\right) \tag{102}
\end{align*}
$$

If $r=s$ then we have

$$
\begin{equation*}
\pi=\left(a_{1}, b_{1}\right) \cdots(k, k+1) \cdots(u, v) \cdots\left(a_{n}, b_{n}\right) \tag{103}
\end{equation*}
$$

- $u<k<k+1<v$

In this case, if $r \neq s$, there is a second partion $\hat{\pi}$ identical to $\pi$ except in the $r^{t h}$ and $s^{\text {th }}$ pairs, $k$ and $k+1$ are switched.

$$
\begin{align*}
& \pi=\left(a_{1}, b_{1}\right) \cdots(u, k) \cdots(k+1, v) \cdots\left(a_{n}, b_{n}\right)  \tag{104}\\
& \hat{\pi}=\left(a_{1}, b_{1}\right) \cdots(u, k+1) \cdots(k, v) \cdots\left(a_{n}, b_{n}\right) \tag{105}
\end{align*}
$$

If $r=s$ then we have

$$
\begin{equation*}
\pi=\left(a_{1}, b_{1}\right) \cdots(u, v) \cdots(k, k+1) \cdots\left(a_{n}, b_{n}\right) \tag{106}
\end{equation*}
$$

- $u<v<k<k+1$

In this case, if $r \neq s$, there is a second partion $\hat{\pi}$ identical to $\pi$ except in the $r^{t h}$ and $s^{t h}$ pairs, $k$ and $k+1$ are switched.

$$
\begin{align*}
& \pi=\left(a_{1}, b_{1}\right) \cdots(u, k) \cdots(v, k+1) \cdots\left(a_{n}, b_{n}\right)  \tag{107}\\
& \hat{\pi}=\left(a_{1}, b_{1}\right) \cdots(u, k+1) \cdots(v, k) \cdots\left(a_{n}, b_{n}\right) \tag{108}
\end{align*}
$$

If $r=s$ then we have

$$
\begin{equation*}
\pi=\left(a_{1}, b_{1}\right) \cdots(u, v) \cdots(k, k+1) \cdots\left(a_{n}, b_{n}\right) \tag{109}
\end{equation*}
$$

Example 3.8. Continuing with the previous example (Example 3.7), with $2 n=8, k=5$ and 2-partition $(1,3)(2,6)(4,8)(5,7)$, the other 2-partition with which this would be paired is $(1,3)(2,5)(4,8)(6,7)$.

Using this construction, we see that after fixing $k$, we may exhaustively list all of the ordered 2-partitions of $\Pi$, identifying each 2-partition as containing a pair $(k, k+1)$ or as being one of the pairs just described.

This allows us to write $P f_{q}$ as a sum of components of the form

$$
\begin{equation*}
(-q)^{*} z_{a_{1}, b_{1}}^{L} \cdots z_{k, k+1}^{L} \cdots z_{a_{n}, b_{n}}^{L} \tag{110}
\end{equation*}
$$

or which appear in pairs such as

$$
\begin{align*}
& (-q)^{*} z_{a_{1}, b_{1}}^{L} \cdots z_{k, u}^{L} z_{k+1, v}^{L} \cdots z_{a_{n}, b_{n}}^{L} \\
& (-q)^{*+1} z_{a_{1}, b_{1}}^{L} \cdots z_{k, v}^{L} z_{k+1, u}^{L} \cdots z_{a_{n}, b_{n}}^{L} \tag{111}
\end{align*}
$$

or

$$
\begin{align*}
& (-q)^{*} z_{a_{1}, b_{1}}^{L} \cdots z_{u, k}^{L} \cdots z_{k+1, v}^{L} \cdots z_{a_{n}, b_{n}}^{L} \\
& (-q)^{*+1} z_{a_{1}, b_{1}}^{L} \cdots z_{u, k+1}^{L} \cdots z_{k, v}^{L} \cdots z_{a_{n}, b_{n}}^{L} \tag{112}
\end{align*}
$$

or

$$
\begin{align*}
& (-q)^{*} z_{a_{1}, b_{1}}^{L} \cdots z_{u, k}^{L} \cdots z_{v, k+1}^{L} \cdots z_{a_{n}, b_{n}}^{L} \\
& (-q)^{*+1} z_{a_{1}, b_{1}}^{L} \cdots z_{u, k+1}^{L} \cdots z_{v, k}^{L} \cdots z_{a_{n}, b_{n}}^{L} \tag{113}
\end{align*}
$$

where $(-q)^{*}$ represents an appropriate power of $(-q)$ determined by $\left(a_{1} b_{1}\right)$ $\left(a_{2} b_{2}\right) \cdots\left(a_{n} b_{n}\right)$. The right action of $e_{k}$ can now be calculated. In the first case, we have the index that contains $(k, k+1)$ and we have

$$
\begin{align*}
q^{*} z_{a_{1}, b_{1}}^{L} & \cdots z_{k, k+1}^{L} \cdots z_{a_{n}, b_{n}}^{L} \cdot e_{k} \\
= & \left(z_{a_{1}, b_{1}}^{L} \cdot e_{k}\right) \cdots\left(z_{k, k+1}^{L} \cdot q^{-\alpha_{k} / 2}\right) \cdots\left(z_{a_{n}, b_{n}}^{L} \cdot q^{-\alpha_{k} / 2}\right) \\
& +\left(z_{a_{1}, b_{1}}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots\left(z_{k, k+1}^{L} \cdot e_{k}\right) \cdots\left(z_{a_{n}, b_{n}}^{L} \cdot q^{-\alpha_{k} / 2}\right) \\
& +\left(z_{a_{1}, b_{1}}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots\left(z_{k, k+1}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots\left(z_{a_{n}, b_{n}}^{L} \cdot e_{k}\right) \\
= & (0) \cdots\left(z_{k, k+1}^{L} \cdot q^{-\alpha_{k} / 2}\right) \cdots\left(z_{a_{n}, b_{n}}^{L} \cdot q^{-\alpha_{k} / 2}\right) \\
& +\left(z_{a_{1}, b_{1}}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots(0) \cdots\left(z_{a_{n}, b_{n}}^{L} \cdot q^{-\alpha_{k} / 2}\right) \\
& +\left(z_{a_{1}, b_{1}}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots\left(z_{k, k+1}^{L} \cdot q^{\alpha_{k} / 2}\right) \cdots(0) \\
= & 0 \tag{114}
\end{align*}
$$

In the next case, with the indexes of the paired 2-partitions containing $(k, u)(k+1, v)$ and $(k, v)(k+1, u)$, the right action of $e_{k}$ can be seen to be zero as well, by using Eq. (92), Eq. (94), and Eq. (95). In fact all remaining
cases are treated similarly, and we get that

$$
\begin{equation*}
P f_{q} \cdot e_{k}=0 \tag{115}
\end{equation*}
$$

A similar argument shows

$$
\begin{equation*}
P f_{q} \cdot f_{k}=0 \tag{116}
\end{equation*}
$$

Since $P f_{q}$ is an element of $A_{q}(X)$ annihilated by the right action of all $e_{k}$ and $f_{k}, 1 \leq k<n, P f_{q}$ must be generated by $\operatorname{det}_{q}$. By comparing degree and coefficients, we see $P f_{q}(Z)=\operatorname{det}_{q}(X)$.

We extend the notation slightly and define

$$
\begin{equation*}
P f_{q}(Z)^{I}=\sum_{\pi \in \Pi^{I}}(-q)^{l(\pi)} z_{\pi}^{L} \tag{117}
\end{equation*}
$$

where $I=\{1,2, \ldots, r\}, r<2 n, r$ is even and $\Pi^{I}$ is the set of ordered 2-partitions of $I$. The proof above also shows that $P f_{q}(Z)^{I}$ is annihilated on the right by all $f_{k}(k<2 n)$ and by all $e_{k}$ except for $r<k<2 n$. As such $P f_{q}(Z)^{I}$ is still a highest weight vector under the right action of $U_{q}(\mathfrak{g})$ and, because it is an element constructed from left symplectic invariant generators, it provides a realization of an element in $A_{q}(X)^{K} \cap A\left(B_{-} \backslash G ; z^{\Lambda_{r}}\right)$.

### 3.5. Decomposition of $\left.{ }^{K} A_{q}(X)\right)$ and $A_{q}(X)^{K}$

We show the decomposition of ${ }^{K} A_{q}(X)$ ) as a right $A(G)$-comodule (resp. left $U_{q}(\mathfrak{g})$-module) and the decomposition of $\left.A_{q}(X)\right)^{K}$ as a left $A(G)$ comodule (resp. right $U_{q}(\mathfrak{g})$-module. To perform this decomposition, several preliminary propositions are presented, along with the introduction of some notational conventions. First some notation:

We define the map $\phi$ from the power set of $\{1,2,3, \ldots, n\}$ into the power set of $\{1,2,3, \ldots, 2 n\}$ by

$$
\begin{equation*}
\phi(A)=\bigcup_{\alpha \in A}\{2 \alpha-1,2 \alpha\} \tag{118}
\end{equation*}
$$

for example $\phi(\{1,3,4,5\})=\{1,2,5,6,7,8,9,10\}$. We will use $\phi$ to construct indices for the rows and columns of quantum minor determinants used in $q$ symplectic invariants and then to describe a specific set of dominant weights
as

$$
\begin{equation*}
P_{2 n}^{\mathcal{A}}=\left\{\lambda \in P_{2 n} ; \lambda=\left(\mu_{1}, \mu_{1}, \mu_{2}, \mu_{2}, \ldots, \mu_{n}, \mu_{n}\right), \mu \in P_{n}\right\} \tag{119}
\end{equation*}
$$

For example $(4,4,4,4,3,3,2,2,2,2,1,1) \in P_{12}^{\mathcal{A}}$.
One of the key ideas used in the decomposition of ${ }^{K} A_{q}(X)$ and $A_{q}(X)^{K}$ is presented in the following proposition (cf. [8]).

Proposition 3.9. Let $\mu \in P_{n}$ be a dominant integral weight and $V_{q}^{R}(\mu)$ be the irreducible left $U_{q}(\mathfrak{g})$ submodule with highest weight $\mu$. Then the space of the $q$-symplectic invariants in $V_{q}^{R}$ has the dimension equal to the multiplicity of $V$ in ${ }^{K} A_{q}(X)$.

Proof. To decompose the algebra $\left.{ }^{K} A_{q}(X)\right)$ as a right $A(G)$-comodule (or left $U_{q}(\mathfrak{g})$ module), it suffices to find the singular weight vectors, i.e. the weight vectors $\phi \in{ }^{K} A_{q}(X)$ such that $e_{k} \cdot \phi=0$ for $k=1, \ldots, n-1$. Since such a singular vector $\varphi$ is contained in the space ${ }^{K} A_{q}(X) \cap A\left(X / B^{+} ; z^{\lambda}\right)$ for some dominant integral weigh $\lambda \in L_{n}$, and generates an irreducible right $A(G)$-comodule with highest weight $\lambda$. Thus if there are $m_{\lambda}$ singular weight vectors of weight $\lambda$ in ${ }^{K} A_{q}(X)$, then the irreducible right $A(G)$-comodule isomorphic to $V_{q}^{R}(\lambda)$ occurs $m_{\lambda}$ times in the decomposition of ${ }^{K} A_{q}(X)$. On the other hand, a singular vector $\varphi$ in ${ }^{K} A_{q}(X) \cap A\left(X / B^{+} ; z^{\lambda}\right)$ is regarded as a left $q$-symplectic invariant in $V_{q}^{L}$ (i.e. annihilated on left). Since $V_{q}^{L}(\lambda)$ and $V_{q}^{R}(\lambda)$ are dual to each other, the dimension of the space of $q$-symplectic invariants coincides.

Next, we show by construction, the existence of a left invariant in the left $U_{q}(\mathfrak{g})$-module $A\left(B_{-} \backslash G ; z^{\lambda}\right)$. We build this left invariant from elements of the following form

$$
\begin{equation*}
a_{r}^{R}=\sum_{J} q^{-2|J|} \xi_{\phi(J)}^{1, \ldots, 2 r} \tag{120}
\end{equation*}
$$

where the sum is over all $J$ such that $\# J=r$ and $J \subseteq\{1,2, \ldots, n\} .|J|$ represents the sum of the elements of $J$. As such $a_{r}^{R} \in A_{q}(X)^{K} \cap A\left(B_{-} \backslash G ; z^{\lambda_{r}}\right)$.

Lemma 3.10. (Existence) For $\lambda=\sum_{r=1}^{n} m_{2 r} \Lambda_{2 r}$, i.e, $\lambda \in P_{2 n}^{\mathcal{A}}, A\left(B_{-} \backslash G ; z^{\lambda}\right)$ contains a left $q$-symplectic invariant.

Proof. Suppose $\lambda=\sum_{r=1}^{n} m_{2 r} \Lambda_{2 r}$, then we define

$$
\begin{equation*}
a_{\lambda}^{R}=\prod_{r=1}^{n}\left(a_{r}^{R}\right)^{m_{2 r}} \tag{121}
\end{equation*}
$$

where $a_{r}^{R}$ is defined by Eq. (120). We see by its construction, $a_{\lambda}^{R} \in A_{q}(X)^{K} \cap$ $A\left(B_{-} \backslash G ; z^{\lambda}\right)$. As such, each right $A(G)$-comodule $V_{q}^{R}(\lambda)$ has a $q$-symplectic invariant.

Lemma 3.11. (Nonexistence) There does not exist a left q-symplectic invariant in the irreducible right $U_{q}(\mathfrak{g})$-submodule $V_{q}^{L}(\lambda)$ if $\lambda \notin P_{2 n}^{\mathcal{A}}$.

Proof. $A_{q}(X)^{K}$ is a right $U_{q}(\mathfrak{g})$-submodule of $A_{q}(X)$. As such, it has its own decomposition into irreducible right $U_{q}(\mathfrak{g})$-submodules indexed by $\lambda \in P_{n}$, where $\lambda$ is a dominant integral weight

$$
\begin{equation*}
A_{q}(X)^{K}=\bigoplus_{\lambda} V_{q}^{L}(\lambda) \tag{122}
\end{equation*}
$$

Each $V_{q}^{L}(\lambda)$ is a highest weight module [16]. Each of these highest weight modules has a realization of $A\left(G / B^{+} ; z^{\lambda}\right)$ with highest weight vector of the form

$$
\begin{equation*}
v_{\lambda}=\left(\xi_{1, \ldots, s}^{1, \ldots, s}\right)^{m_{s}}\left(\xi_{1, \ldots, s-1}^{1, \ldots, s-1}\right)^{m_{s-1}} \cdots\left(\xi_{1}^{1}\right)^{m_{1}} \tag{123}
\end{equation*}
$$

However, because, the elements of $A_{q}(X)^{K}$ are annihilated on the left by all $e_{k}$ and $f_{k}$ where $k$ is odd, then for the highest weight vector $v_{\lambda}$, it must be true that $\lambda=\sum_{r=1}^{n} m_{r} \Lambda_{r}$ where $m_{r}=0$ when $r$ is odd. Thus,

$$
\begin{equation*}
A_{q}(X)^{K}=\bigoplus_{\lambda \in P_{2 n}^{\mathcal{A}}} V_{q}^{L}(\lambda) \tag{124}
\end{equation*}
$$

Lemma 3.12. (Uniqueness) The multiplicity of $V_{q}^{R}(\lambda)$ an irreducible right $U_{q}(\mathfrak{g})$-module, in the decomposition of $A_{q}(X)^{K}$ is exactly one.

Proof. As mentioned earlier, by Proposition 3.9, the multiplicity of $V_{q}^{R}(\lambda)$ in the decomposition of $A_{q}(X)^{K}$ is equal to the number of left $q$-symplectic invariants in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$. Let $v^{K}$ be a non zero left invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$,
as such, it can be written as a linear combination of weight vectors from the standard basis of $A\left(B_{-} \backslash G ; z^{\lambda}\right)[16]$.

However, since $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ is a highest weight vector space, there must be at least one basis (weight) vector, $\eta$, in the composition of $v^{K}$ for which there are no higher weight vectors in $v^{K}$. In other words

$$
\begin{equation*}
v^{K}=\eta \oplus v_{1} \oplus \cdots v_{j} \tag{125}
\end{equation*}
$$

where the weights of $v_{1}, \ldots, v_{j}$ are less than or equal to that of $\eta$. As such, $\eta$ must be annihilated by all $e_{k}$, where $k<2 n$ and $k$ is odd. Additionally, the elements of the form

$$
\begin{align*}
s p_{f}(i, i+1)= & e_{2 i}+q^{-2}\left(f_{2 i-1} f_{2 i} f_{2 i+1}-f_{2 i} f_{2 i-1} f_{2 i+1}\right. \\
& \left.-f_{2 i+1} f_{2 i-1} f_{2 i}+f_{2 i+1} f_{2 i} f_{2 i-1}\right) \tag{126}
\end{align*}
$$

where $1 \leq i<n$, also annihilate $v^{K}$, and this in turn requires that $\eta$ also be annihilated by all $e_{k}$, where $k<2 n$ and $k$ is even. Therefore $\eta$ must be a highest weight vector of $A\left(B_{-} \backslash G ; z^{\lambda}\right)$, but this vector is unique up to constant multiple, because $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ is a highest weight module. So

$$
\begin{equation*}
\eta=c v_{\lambda}, \quad c \in \mathbb{C} \tag{127}
\end{equation*}
$$

where $v_{\lambda}$ is defined by Eq. (123). This tells us that any non-zero left $q$ symplectic invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ must be written as

$$
\begin{equation*}
c v_{\lambda} \oplus w_{1} \oplus \cdots w_{j}, \quad c \in \mathbb{C}, c \neq 0 \tag{128}
\end{equation*}
$$

where $w_{1}, \ldots, w_{j}$ are lower weight vectors of $A\left(B_{-} \backslash G ; z^{\lambda}\right)$.
Now assume there is more than one left quantum $q$-symplectic invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$, say $v^{K}$ and $w^{K}$. Each of these may be written as a sum of standard basis elements, each including a non-zero term for the highest weight vector $v_{\lambda}$. In other words, they may be written as

$$
\begin{align*}
v^{K} & =c_{0} v_{\lambda}+c_{1} v_{1}+c_{2} v_{2}+\cdots+c_{i} v_{i}, \quad c_{0} \neq 0  \tag{129}\\
w^{K} & =k_{0} v_{\lambda}+d_{1} v_{1}+d_{2} v_{2}+\cdots+d_{j} v_{j}, \quad k_{0} \neq 0 \tag{130}
\end{align*}
$$

Since the linear combination of any left $q$-symplectic invariant is also a left $q$-symplectic invariant then it must be true that $k_{0} v^{K}-c_{0} w^{K}$ is also
a left $q$-symplectic invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$. If $k_{0} v^{K}-c_{0} w^{K} \neq 0$ then we have a contradiction to the requirement that any left $q$-symplectic invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ has a nonzero $v_{\lambda}$ component. On the other hand, if $k_{0} v^{K}-c_{0} w^{K}=0$ then $w^{K}$ is a constant multiple of $v^{K}$. Therefore, any left $q$-symplectic invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ is unique up to a constant multiple.

The following proposition summarizes Lemmas 3.10, 3.11, and 3.12,
Proposition 3.13. The space of $q$-symplectic invariants in the right $A(G)$ comodule $V_{q}^{R}(\mu)$ is one dimensional if and only if $\mu=\sum_{r=1}^{n} m_{2 r} \Lambda_{2 r}$, in other words, $\mu \in P_{2 n}^{\mathcal{A}}$. Otherwise there are no $q$-symplectic invariants in $V_{q}^{R}$

By Proposition 3.9 we may then summarize our results with the following theorem

Theorem 3.14. The irreducible decomposition of $A_{q}(X)^{K}$ as a right $U_{q}(\mathfrak{g})$ module is given by

$$
\begin{equation*}
A_{q}(X)^{K}=\bigoplus_{\lambda \in P_{2 n}^{\mathcal{A}}} V_{q}^{L}(\lambda) \tag{131}
\end{equation*}
$$

similarly ${ }^{K} A_{q}(X)$, as a left $U_{q}(\mathfrak{g})$-module has the irreducible decomposition

$$
\begin{equation*}
{ }^{K} A_{q}(X)=\bigoplus_{\lambda \in P_{2 n}^{\mathcal{A}}} V_{q}^{R}(\lambda) \tag{132}
\end{equation*}
$$

Where $P_{2 n}^{\mathcal{A}}$ is defined by Eq. (119).
Proposition 3.15. The space $A_{q}^{L}(\mathcal{A})=A_{q}(X)^{K}$, (resp. $\left.{ }^{K} A_{q}(X)=A_{q}^{R}(\mathcal{A})\right)$. As such, $A_{q}^{L}(\mathcal{A})\left(\right.$ resp. $\left.A_{q}^{R}(\mathcal{A})\right)$ also have the decompositions as a right (resp. left) $U_{q}(\mathfrak{g})$-modules,

$$
\begin{align*}
& A_{q}^{L}(\mathcal{A})=\bigoplus_{\lambda \in P_{2 n}^{\mathcal{A}}} V_{q}^{L}(\lambda)  \tag{133}\\
& A_{q}^{R}(\mathcal{A})=\bigoplus_{\lambda \in P_{2 n}^{\mathcal{A}}} V_{q}^{R}(\lambda) \tag{134}
\end{align*}
$$

Proof. From its definition, we already have $A_{q}^{L}(\mathcal{A}) \subseteq A_{q}(X)^{K}$. The elements, $P f_{q}^{I}$ described by Eq. (117) provide a formula for explicitly constructing a left $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$ invariant in $A\left(B_{-} \backslash G ; z^{\lambda}\right)$ for any $\lambda$. As such, $A_{q}(X)^{K} \subseteq A_{q}^{L}(\mathcal{A})$, and we have $A_{q}^{L}(\mathcal{A})=A_{q}(X)^{K}$.

## 4. Bi-invariants

In this section we define a subalgebra of $A_{q}(X)$ by the intersection of $A_{q}^{R}(\mathcal{A})$ and $A_{q}^{L}(\mathcal{A})$. Defined in this way, this space is annihilated on the left and right by $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$. We then proceed to show that this algebra is really $\mathbb{C}\left[s_{1}, \ldots, s_{n}\right]^{\mathfrak{G}_{n}}$, the symmetric algebra of $n$ variables. To start, we define $A_{Z P}$, as

$$
\begin{equation*}
A_{Z P}=A_{q}^{R}(\mathcal{A}) \cap A_{q}^{L}(\mathcal{A})=\bigoplus_{m=0}^{\infty} A_{Z P, 2 m}, \tag{135}
\end{equation*}
$$

Recall, the polynomials of $A_{q}^{R}(\mathcal{A})$ and $A_{q}^{L}(\mathcal{A})$ have even degree so it has the natural grading into the subspaces $A_{Z P, 2 m}$.

Now we define

$$
\begin{equation*}
E_{r}=\sum_{I, J} q^{2(|I|-|J|)} \xi_{\phi(J)}^{\phi(I)}, \quad 1 \leq r \leq n \tag{136}
\end{equation*}
$$

where the summation runs over all subsets $I$ and $J$ of $\{1, \ldots, n\}$ and $\# I=$ $\# J=r$. Here, $|I|$ and $|J|$ are the sums of the elements of $I$ and $J$ respectively.

Lemma 4.16. $E_{r} \in A_{Z P, 2 r}$

Proof. If we examine the component of $E_{r}$ that is obtained by holding $I$ fixed at $I=\{1,2, \ldots, r\}$, we see that this component is precisely $a_{r}^{R}$, defined in Eq. (120). As such, this component is invariant under the left action of $U_{q}(\mathfrak{s p}(2 n, \mathbb{C}))$. The remaining components of $E_{r}$ (the components obtained by fixing $I$ at other values) can be obtained by the right action of $U_{q}(\mathfrak{g})$ on $a_{r}^{R}$. Since $A_{q}(X)^{K}$ is a right submodule of $A_{q}(X)$ these other components of $E_{r}$ must also be left invariant. Thus, $E_{r} \in A_{q}(X)^{K}$. Similarly, we see that the component of $E_{r}$ associated with the fixed $J=\{1,2, \ldots, n\}$ is in ${ }^{K} A_{q}(X)$ and likewise the other components of $E_{r}$ can be obtained by the
left action of $U_{q}(\mathfrak{g})$. Thus, $E_{r} \in{ }^{K} A_{q}(X)$. Since $E_{r}$ has degree $2 r$ (by its construction) and $E_{r} \in A_{q}^{R}(\mathcal{A}) \cap A_{q}^{L}(\mathcal{A})$, it follows that $E_{r} \in A_{Z P, 2 r}$.

Theorem 4.17. The algebra $A_{Z P}$ is generated by $E_{r}(1 \leq r \leq n)$ and the algebra $A_{Z P}$ is isomorphic to the algebra of symmetric polynomials of $n$ variables;

$$
\begin{equation*}
\pi: A_{Z P} \xrightarrow[\rightarrow]{\tilde{C}}\left[s_{1}, \ldots, s_{n}\right]^{\mathfrak{S}_{n}} \tag{137}
\end{equation*}
$$

Proof. Because of the decomposition given in Proposition 3.15, the dimension of the bi-invariant space associated with each $\lambda \in P_{2 n}^{\mathcal{A}}$ must be exactly one. Since the degree of the polynomial in each of these bi-invariant spaces is $\sum_{k=1}^{n} \lambda_{k}$, the dimension of $A_{Z P, 2 m}$ can then be calculated as the number of partitions in $P_{2 n}^{\mathcal{A}}$ of $2 m$. As these partitions are in $P_{2 n}^{\mathcal{A}}$ we may also consider this as the number of partitions of $m$ whose number of parts is less than or equal to $n$. Adopting the notation of Jing and Yamada [8] we denote this by $p_{n}(m)$.

Consider the restriction of the projection map $\pi$ to $A_{Z P}$

$$
\begin{equation*}
\pi_{H}^{\prime}: A_{Z P} \rightarrow A_{+}(H) \tag{138}
\end{equation*}
$$

where $A_{+}(H)=\mathbb{C}\left[t_{1}, \ldots, t_{n}\right]$. Then $\operatorname{Ker}\left(\pi_{H}^{\prime}\right)=\bigoplus_{r=0}^{\infty} \operatorname{Ker}\left(\pi_{H, 2 r}^{\prime}\right)$, where

$$
\begin{equation*}
\pi_{H, 2 r}^{\prime}: A_{Z P, 2 r} \rightarrow A_{2 r}(H) \tag{139}
\end{equation*}
$$

Similar to the proof by [8], the monomials $E_{r_{1}} E_{r_{2}} \ldots E_{r_{k}}\left(r_{1} \leq r_{2} \leq \ldots \leq r_{k}\right)$ have the degree $2\left(r_{1}+r_{2}+\ldots+r_{k}\right)$ and are linearly independent over $\mathbb{C}$. As such the space of degree $2 m$ spanned by these monomials has dimension $p_{n}(m)$. This shows that the space $A_{Z P}$ is generated by $E_{r}(1 \leq r \leq n)$.

Additionally, the map $\pi_{H, 2 r}^{\prime}$ acts on the generators of $A_{Z P}$ in the following manner

$$
\begin{aligned}
\pi^{\prime}\left(E_{r}\right) & =\pi^{\prime}\left(\sum_{I} \xi_{\phi(I)}^{\phi(I)}\right) \\
& =\sum_{I}\left(t_{2 i_{1}-1} t_{2 i_{1}}\right)\left(t_{2 i_{2}-1} t_{2 i_{2}}\right) \cdots\left(t_{2 i_{r}-1} t_{2 i_{r}}\right) \neq 0
\end{aligned}
$$

where the sum runs over all subsets $I$ of $\{1,2, \ldots, n\}$ and $\# I=r$, thus $\operatorname{Ker}\left(\pi_{H, 2 r}^{\prime}\right)=(0)$. Another way of viewing this is that each of these $E_{r}$ has monomials which are products of diagonal elements. As such, $\pi\left(E_{r}\right) \neq 0$ for $1 \leq r \leq n$. Thus we have the isomorphism

$$
\begin{aligned}
A_{Z P} & \cong \mathbb{C}\left[\left(t_{1} t_{2}\right),\left(t_{3} t_{4}\right), \ldots,\left(t_{n-1} t_{n}\right)\right]^{\mathfrak{S}_{n}} \\
& \cong \mathbb{C}\left[s_{1}, s_{2}, \ldots, s_{n}\right]^{\mathfrak{S}_{n}}
\end{aligned}
$$

where we let $s_{i}=t_{2 i-1} t_{2 i}$.

## 5. Spherical Functions and Symmetric Polynomials

Through the isomorphism in Theorem (4.1) our $q$-zonal polynomials are basis elements in the ring of symmetric polynomials, and they are clearly $q$-deformations of the zonal polynomials defined on $G L(2 n, \mathbb{C}) / S p(2 n, \mathbb{C})$. We describe the relation with Macdonald polynomials [13].

Macdonald polynomials are special orthogonal basis of the commutative algebra
$\mathbb{Q}(q, t)\left[x_{1}, \ldots, x_{n}\right]^{\mathfrak{G}_{n}}$, where $q$ and $t$ are two parameters. To describe them we consider the following shift operator $T_{u, x_{i}}$ by

$$
\left(T_{u, x_{i}} f\right)\left(x_{1}, \ldots, x_{n}\right)=f\left(x_{1}, \ldots, u x_{i}, \ldots, x_{n}\right)
$$

for each $f \in \mathbb{Q}(q, t)\left[x_{1}, \ldots, x_{n}\right]$. Let $X$ be another indeterminate and define

$$
\begin{aligned}
D(X, q, t) & =\Delta^{-1} \sum_{w \in \mathfrak{S}_{n}} \epsilon(w) z^{w \delta} \prod_{i=1}^{n}\left(X+t^{(w \delta)_{i}} T_{q, x_{i}}\right) \\
& =\sum_{r=0}^{n} D_{r} X^{n-r}
\end{aligned}
$$

where $\delta=(n-1, n-2, \ldots, 1,0)$ and

$$
\Delta=\prod_{1 \leq i<j \leq n}\left(x_{i}-x_{j}\right)
$$

is the Vandermonde determinant in $x_{1}, \ldots, x_{n}$. It follows immediately that
$D_{0}=1$ and

$$
D_{1}=\sum_{i=1}^{n}\left(\prod_{j \neq i} \frac{t x_{i}-x_{j}}{x_{i}-x_{j}}\right) T_{q, x_{i}}
$$

Macdonald showed that for each partition $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$ there is a unique symmetric polynomial $P_{\lambda}(x ; q, t)$ satisfying the two conditions (4.54.6):

$$
\begin{equation*}
P_{\lambda}=m_{\lambda}+\sum_{\mu<\lambda} u_{\lambda \mu} m_{\mu} \tag{140}
\end{equation*}
$$

where $u_{\lambda \mu} \in \mathbb{Q}(q, t)$ and $m_{\mu}=x_{1}^{\mu_{1}} \ldots x_{n}^{\mu_{n}}+\ldots$ is the monomial symmetric polynomial;

$$
\begin{equation*}
D_{1} P_{\lambda}=\left(\sum_{i=1}^{n} q^{\lambda_{i}} t^{n-i}\right) P_{\lambda} \tag{141}
\end{equation*}
$$

Moreover Macdonald proves that $P_{\lambda}$ is also an eigenfunction for all the difference operators $D_{r}$, and

$$
\begin{equation*}
D(X ; q, t) P_{\lambda}=\prod_{i=1}^{n}\left(X+t^{n-i} q^{\lambda_{i}}\right) P_{\lambda} \tag{142}
\end{equation*}
$$

The polynomial $P_{\lambda}(x ; q, t)$ is called the Macdonald polynomial associated with the partition $\lambda$. In particular, $P_{\lambda}(x ; q, q)$ is the famous Schur polynomial; $\lim _{t \rightarrow 1} P_{\lambda}\left(x ; t^{2}, t\right)$ is the zonal polynomial.

Proposition 5.18. Under the isomorphism $\pi: A_{z p} \longrightarrow \mathbb{C}\left[z_{1}, \ldots, z_{n}\right]^{\mathfrak{G}_{n}}$, the $q$-zonal polynomial in $V_{q}(\lambda)$ is a constant multiple of the Macdonald polynomial $P_{\lambda}\left(z ; q^{2}, q^{-4}\right)$.

The general case of quantum spherical functions was studied by Noumi [15] using quantum groups and Letzter [12] using quantum enveloping algebras. In the following we will outline a different approach to understand the relationship between symmetric functions and quantum invariants. First of all let's study the $q$-difference operators on $V_{q}(2 \lambda)$.

Recall the center of the quantized universal enveloping algebra $U_{q}\left(\mathfrak{s l} l_{n-1}\right)$
is generated by the following $n-1$ elements [17].

$$
c_{k}=\sum_{\sigma, \sigma^{\prime} \in \mathfrak{S}_{n}}(-q)^{l(\sigma)+l\left(\sigma^{\prime}\right)} l_{\sigma_{1}, \sigma_{1}^{\prime}}^{(+)} \cdots l_{\sigma_{k} \sigma_{k}^{\prime}}^{(+)} \sigma_{\sigma_{k+1} \sigma_{k+1}^{\prime}}^{(-)} \cdots l_{\sigma_{n} \sigma_{n}^{\prime}}^{(-)}, \quad k=1, \cdots, n-1
$$

where $L^{( \pm)}=\left(l_{i j}^{( \pm)}\right)$is the upper (lower) triangular defining matrix for the quantum algebra $U_{q}\left(s l_{n-1}\right)$ in the FRT formulation [17] and $l(\sigma)=\#\{i<$ $\left.j \mid \sigma_{i}>\sigma_{j}\right\}$. We only remark that the elements $l_{i j}^{( \pm)}$are analogs of Weylgenerators for $U_{q}\left(\mathfrak{s l}_{n-1}\right)$. In particular

$$
l_{i i}^{( \pm)}=q^{ \pm \epsilon_{i}}
$$

The algebra $U_{q}\left(\mathfrak{s l} l_{n-1}\right)$ acts on $G L_{q}(n, \mathbb{C})$ as q-difference operators, thus the center of $U_{q}\left(\mathfrak{s l}_{n-1}\right)$ acts on modules $V_{q}(2 \lambda)$ as scalar operators. In particular our $q$-zonal polynomials are simultaneous eigenfunctions of these $q$ difference operators.

Theorem 5.19. For $1 \leq k \leq n-1$, the central element $c_{k}$ acts on the irreducible $U_{q}\left(s l_{n}\right)$-module $V(\lambda)$ as a scalar multiplication by

$$
q^{2|\lambda|+\binom{n}{2}+k(n-1)}[k]![n-k]!\left(\sum_{1 \leq i_{1}<\cdots<i_{k} \leq n} q^{-2 \lambda_{i_{1}}-\cdots-2 \lambda_{i_{k}}+2\left(i_{1}-n\right)+\cdots+2\left(i_{k}-n\right)}\right),
$$

where $|\lambda|=\lambda_{1}+\cdots+\lambda_{n}$.

Proof. Pick a lowest weight vector $v_{0}$ in $V(\lambda)$ with the weight $-\lambda=-\lambda_{1} \epsilon_{1}-$ $\cdots-\lambda_{n} \epsilon_{n}$. Note that the generators $l_{i j}^{(+)}, l_{j i}^{(-)}(i<j)$ belong to the so-called strict upper and lower Borel subalgebra generated by $e_{i}$ and $f_{i}(i=1, \cdots n-$ 1) respectively. The element $l_{\sigma_{k+1} \sigma_{k+1}^{\prime}}^{(-)} \cdots l_{\sigma_{n} \sigma_{n}^{\prime}}^{(-)}$kills $v_{0}$ unless $\sigma_{k+1}=\sigma_{k+1}^{\prime}$, $\ldots, \sigma_{n}=\sigma_{n}^{\prime}$. But $\sigma_{1} \leq \sigma_{1}^{\prime}, \ldots, \sigma_{k} \leq \sigma_{k}^{\prime}$, so one must have $\sigma=\sigma^{\prime}$ in the action of $c_{n-k}$ on $v_{0}$. We thus have

$$
\begin{aligned}
c_{k} v_{0} & =\sum_{\sigma \in \mathfrak{S}_{n}} q^{2 l(\sigma)} q^{-\lambda_{\sigma_{1}}-\cdots-\lambda_{\sigma_{k}}+\lambda_{\sigma_{k+1}}+\cdots+\lambda_{\sigma_{n}}} v_{0} \\
& =q^{|\lambda|} \sum_{\sigma \in \mathfrak{S}_{n}} q^{2 l(\sigma)-2 \lambda_{\sigma_{1}}-\cdots-2 \lambda_{\sigma_{k}}} v_{0}
\end{aligned}
$$

Consider the Young subgroup $\mathfrak{S}_{k} \times \mathfrak{S}_{n-k}$ of $\mathfrak{S}_{n}$. We can choose its left coset representatives to be the elements $\tau$ such that $\tau_{1}<\cdots<\tau_{k}$,
$\tau_{k+1}<\cdots<\tau_{n}$. Recall that an inversion of the permutation $\tau$ is a pair $(i j)$ such that $i<j$ and $\tau_{i}>\tau_{j}$. By construction the inversions of $\tau$ may only take place among $(i j)$ where $i \leq k$ and $j \geq k+1$. For each $i(\leq k)$, there are $\tau_{i}-1$ natural numbers less than $\tau_{i}$, and $i-1$ of them already appear before $\tau_{i}$ in the permutation. So there are $\tau_{i}-i$ inversions of $\tau$ in the form $(i j)$, which implies that $l(\tau)=\sum_{i=1}^{k}\left(\tau_{i}-i\right)$.

Let $\tau \sigma$ be the general element in $\mathfrak{S}_{n}$ where $\sigma=\sigma_{1} \sigma_{2} \in \mathfrak{S}_{k} \times \mathfrak{S}_{n-k}$. In the sequence $(\tau \sigma(1), \ldots, \tau \sigma(k), \tau \sigma(k+1), \ldots, \tau \sigma(n))$ we divide the inversions of $\tau \sigma$ into three parts: the inversions among the first $k$ numbers, those among the last $n-k$ numbers, and the inversions between the first $k$ numbers and the last $n-k$ numbers. The second part $(\tau \sigma(k+1), \ldots, \tau \sigma(n))=\left(\tau \sigma_{2}(k+\right.$ $\left.1), \ldots, \tau \sigma_{2}(n)\right)$ has $l\left(\sigma_{2}\right)$ inversions as $\tau$ preserves the order of $k+1, \ldots, n$, similarly the first part $(\tau \sigma(1), \ldots, \tau \sigma(k))=\left(\tau \sigma_{1}(1), \ldots, \tau \sigma_{1}(k)\right)$ has $l\left(\sigma_{1}\right)$ inversions among them. Observe that we are free to switch the numbers in each part when considering the inversions between the first part and the second part, thus the number of inversions of this type are exactly $l(\tau)$. Therefore we have

$$
\begin{aligned}
l\left(\tau \sigma_{1} \sigma_{2}\right) & =l(\tau)+l\left(\sigma_{1}\right)+l\left(\sigma_{2}\right) \\
& =l\left(\sigma_{1}\right)+l\left(\sigma_{2}\right)+\sum_{i=1}^{k}\left(\tau_{i}-i\right)
\end{aligned}
$$

where $\sigma_{1} \in \mathfrak{S}_{k}, \sigma_{2} \in \mathfrak{S}_{n-k}, \tau \in \mathfrak{S}_{n} /\left(\mathfrak{S}_{k} \times \mathfrak{S}_{n-k}\right)$.
Now let's return back to the action $c_{k} v_{o}$. Using the invariance of $\lambda_{\sigma(1)}+$ $\cdots+\lambda_{\sigma(k)}$ under $\mathfrak{S}_{k} \times \mathfrak{S}_{n-k}$, we have that

$$
\begin{aligned}
c_{k} v_{0}= & q^{2|\lambda|} \sum_{\tau, \sigma_{1}, \sigma_{2}} q^{2 l\left(\tau \sigma_{1} \sigma_{2}\right)-2 \lambda_{\tau \sigma_{1} \sigma_{2}(1)}-\cdots-2 \lambda_{\tau \sigma_{1} \sigma_{2}(k)}} v_{0} \\
= & q^{2|\lambda|} \sum_{\tau, \sigma_{1}, \sigma_{2}} q^{2 l\left(\tau \sigma_{1} \sigma_{2}\right)-2 \lambda_{\tau(1)}-\cdots-2 \lambda_{\tau(k)}} v_{0} \\
= & q^{2|\lambda|} \sum_{\sigma_{1} \in \mathfrak{S}_{k}} q^{2 l\left(\sigma_{1}\right)} \sum_{\sigma_{2} \in \mathfrak{S}_{n-k}} q^{2 l\left(\sigma_{2}\right)} \sum_{\tau} q^{2 l(\tau)-2 \lambda_{\tau(1)}-\cdots-2 \lambda_{\tau(k)}} v_{0} \\
= & q^{2|\lambda|+\binom{k}{2}+\binom{n-k}{2}}[k]![n-k]! \\
& \cdot\left(\sum_{\tau(1)<\cdots<\tau(k)} q^{-2 \lambda_{\tau(1)}-\cdots-2 \lambda_{\tau(k)}+2(\tau(1)-1)+\cdots+2(\tau(k)-k)}\right) v_{0} \\
= & q^{2|\lambda|+\binom{n}{2}+k(n-1)}[k]![n-k]!
\end{aligned}
$$

$$
\cdot\left(\sum_{1 \leq \tau(1)<\cdots<\tau(k) \leq n} q^{-2 \lambda_{\tau(1)}-\cdots-2 \lambda_{\tau(k)}+2(\tau(1)-n)+\cdots+2(\tau(k)-n)}\right) v_{0}
$$

where we have used the well-known identity $\sum_{\sigma \in \mathfrak{S}_{n}} q^{2 l(\sigma)}=q^{\binom{n}{2}}[n]!(\operatorname{cf} .[2])$.

Now we restrict ourselves to the case of irreducible highest $U_{q}\left(s l_{2 n}\right)$ module $V(\tilde{\lambda})$ such that $\tilde{\lambda}=\tilde{\lambda}_{1} \epsilon_{1}+\cdots+\tilde{\lambda}_{2 n} \epsilon_{2 n}$ and $\lambda_{2 i-1}=\lambda_{2 i}=\lambda_{i}$ for $i=1, \ldots, n$. It is also a lowest weight module with the lowest weight $-\tilde{\lambda}$.

Theorem 5.20. The bi-invariant function inside $V(\tilde{\lambda})$, restricted to the ring of symmetric functions, is the Macdonald symmetric function $P_{\lambda}\left(q^{2}, q^{4}\right)$.

Proof. It follows from the theorem in the case of $U_{q}\left(s l_{2 n}\right)$-module $V(\tilde{\lambda})$ that

$$
\begin{aligned}
c_{1} v_{0} & =q^{2|\tilde{\lambda}|+\binom{2 n}{2}+2(2 n-1)}[2]![2 n-2]!\sum_{1 \leq i \leq 2 n} q^{-2 \tilde{\lambda}_{i}+2(i-2 n)} v_{0} \\
& =q^{4|\lambda|+\binom{2 n}{2}+2(2 n-1)-1}[2]^{2}[2 n-2]!\sum_{1 \leq i \leq n} q^{-2 \lambda_{i}+4(i-n)} v_{0} .
\end{aligned}
$$

In other words, the quantum Casimir operator $c_{1}$ agrees with Macdonald operator $D_{1}\left(q^{2}, q^{4}\right)$ or $D_{1}\left(q^{-2}, q^{-4}\right)$ on the space. We note that the leading term of the spherical functions, when restricted to the zonal part, is exactly the leading term of the Macdonald spherical function $P_{\lambda}\left(q^{2}, q^{4}\right)$ (which also agrees with Schur function $s_{\lambda}$ ). Hence the eigenfunction restricted to the ring of symmetric functions is the Macdonald symmetric function $P_{\lambda}\left(q^{2}, q^{4}\right)$. Similarly the action of the higher difference operators are given by

$$
\begin{aligned}
c_{2 k} v_{0} & =q^{2|\tilde{\lambda}|+\binom{2 n}{2}+2 k(2 n-1)}[2 k]![2 n-2 k]! \\
& \cdot\left(\sum_{1 \leq \tau(1)<\cdots<\tau(2 k) \leq n} q^{-2 \tilde{\lambda}_{\tau(1)}-\cdots-2 \tilde{\lambda}_{\tau(2 k)}+2(\tau(1)-2 n)+\cdots+2(\tau(2 k)-2 n)}\right) v_{0}
\end{aligned}
$$

The last identity plus the same idea also gives that
Corollary 5.21. The restriction of $c_{k}$ to the ring $\mathbb{C}\left[z_{1}, \ldots, z_{n}\right]^{\mathfrak{S}_{n}}$ is exactly the difference operator $D_{k}\left(q^{2}, q^{4}\right), k=1, \ldots, n$ up to a constant.

## Acknowledgment

NJ gratefully acknowledges the partial support of Max-Planck Institut für Mathematik in Bonn, Simons Foundation grant 198129, and NSFC grant 10728102 during this work.

## References

1. E. Abe, Hopf Algebras. Cambridge Tracts in Mathematics, 74. Cambridge University Press, Cambridge-New York, 1980.
2. N. Bourbaki, Groupes et algébres de Lie, Ch. 4-6, Mason, Paris, 1981.
3. R. Howe, Remarks on classical invariant theory, Trans. Amer. Math. Soc. 313 (1989), no. 2, 539-570.
4. L. K. Hua, Harmonic analysis of functions of several complex variables in the classical domains, AMS Translations 6, Providence, RI, 1963.
5. N. Jacobson, Basic Algebra, 2nd ed., W. H. Freeman and Company, 1985.
6. A. T. James, Zonal polynomials of the real positive definite symmetric matrices, Ann. Math., 74 (1961), 456-469.
7. M. Jimbo, A $q$-difference analogue of $U(g)$ and the Yang-Baxter equation, Lett. Math. Phys., 10 (1985), 63-69.
8. N. Jing and H.-F. Yamada, Zonal polynomials on the quantum general linear groups, In: Nankai Workshop on Quantum groups, Edited by M. L. Ge, World Sci, Singapore, 1995, 66-72.
9. C. Kassel, Quantum Groups, Springer-Verlag, New York, 1995.
10. T. Kornwinder, Orthogonal polynomials in connections with quantum groups, In: Orthogonal Polynomials, Edited by P. Nevai, NATO ASI ser., Kluwer Acad. Publishers, 1990, 257-297.
11. G. I. Lehrer, H. Zhang and R. Zhang, A quantum analogue of the first fundamental theorem of classical invariant theory, Comm. Math. Phys., 301 (2011), No. 1, 131-174.
12. G. Letzter, Quantum zonal spherical functions and Macdonald polynomials, Adv. Math., 189 (2004), No. 1, 88-147.
13. I. G. Macdonald, Symmetric Functions and Hall Polynomials, 2nd ed., Clarendon Press, Oxford, 1995.
14. S. Montgomery, Hopf algebras and their actions on rings, CBMS ser. 82, AMS, Providence, RI., 1993.
15. M. Noumi, Macdonald's symmetric polynomials as zonal symmetric functions on some quantum homogeneous spaces, Adv. Math., 123(1996), 16-77.
16. M. Noumi, H. Yamada and K. Mimachi, Finite dimensional representations of the quantum groups $g l_{q}(n ; C)$ and the zonal spherical functions on $u_{q}(n-1) \backslash U_{q}(n)$, Japanese J. of Math., (1993), 31-80.
17. N. Yu. Reshetikhin, L. A. Takhtajan and L. D. Faddeev, Quantization of Lie groups and Lie algebras, Algebra and Analysis 1 (1989), 178-206; English Transl. Leningrad Math. J. 1(1990), 193-225.
18. E. Strickland, Classical invariant theory for the quantum symplectic group. Adv. Math., 123 (1996), No. 1, 78-90.
19. K. Ueno and T. Takebayashi, Zonal spherical functions on quantum symmetric spaces and Macdonald's symmetric polynomials, In Quantum Groups, Edited by P. Kulish, Lect. Notes Math. vol. 1510, Springer-Verlag, 1992, 142-147.

[^0]:    Received September 26, 2011 and in revised form October 6, 2011.
    AMS Subject Classification: Primary: 20G42; Secondary: 17B37, 43A90, 05E10.
    Key words and phrases: Pfaffians, quantum groups, invariants, quantum anti-symmetric matrices.

