# QUANTITATIVE ESTIMATE OF THE STATIONARY NAVIER-STOKES EQUATIONS AT INFINITY AND UNIQUENESS OF THE SOLUTION

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This work is dedicated to Professor Tai-Ping Liu for his 70th birthday.

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#### Abstract

In this paper we are interested in the asymptotic behavior of incompressible fluid around a bounded obstacle. Under certain a priori decaying assumptions, we derive a quantitative estimate of the decaying rate of the difference of any two velocity functions at infinity. This quantitative estimate gives us a sufficient condition, expressed in terms of integrability, to guarantee that the solution of the Navier-Stokes equations is unique.

#### 1. Introduction

Let D be a bounded domain in  $\mathbb{R}^n$  and  $\Omega = \mathbb{R}^n \setminus \overline{D}$  with  $n \geq 2$ . Without loss of generality, we let 0 belong to interior of D. Assume that  $\Omega$  is filled with an incompressible fluid described by the stationary Navier-Stokes equations

$$\begin{cases} -\Delta u + u \cdot \nabla u + \nabla p = f & \text{in } \Omega, \\ \nabla \cdot u = 0 & \text{in } \Omega. \end{cases}$$
(1.1)

We are interested in the following question: let  $u_1$  and  $u_2$  be two solutions of (1.1) satisfying some pre-described assumptions such as boundedness or

AMS Subject Classification: 35B60, 76D05.

Received March 31, 2015 and in revised form September 2, 2015.

Key words and phrases: Quantitative uniqueness estimates, Navier-Stokes equations, Carleman estimates.

decaying conditions, then find a sufficient condition which guarantees that  $u_1 \equiv u_2$  in  $\Omega$ . In this paper, we answer this question by deriving a minimal decay rate of  $u_1 - u_2$  at infinity if  $u_1 \neq u_2$ .

This question is motivated by the following problem. It was shown by Finn [1] that when n = 3 and f = 0, if  $u|_{\partial D} = 0$  and  $u = o(|x|^{-1})$ , then u is trivial. Inspired by Finn's result, we would like to ask the following question: when n = 3, if we know a priori that  $u = O(|x|^{-1})$ , what is the minimal decaying rate of any nontrivial u satisfying (1.1)? It should be remarked that the boundary value of u on  $\partial B$  is irrelevant in this problem. Moreover, the asymptotic behavior  $u = O(|x|^{-1})$  characterizes the so-called physically reasonable solutions introduced by Finn [2].

To answer the main question of the paper, we simply subtract two equations for  $u_1$  and  $u_2$  and obtain

$$\begin{cases} -\Delta v + v \cdot \nabla v + v \cdot \nabla u_2 + u_2 \cdot \nabla v + \nabla p_v = 0 & \text{in } \Omega, \\ \nabla \cdot v = 0 & \text{in } \Omega, \end{cases}$$

where  $v = u_1 - u_2$  and  $p_v = p_1 - p_2$ . Therefore, to solve the problem, it suffices to consider the generalized Navier-Stokes equations

$$\begin{cases} -\Delta v + v \cdot \nabla v + v \cdot \nabla \alpha + \alpha \cdot \nabla v + \nabla p = 0 & \text{in } \Omega, \\ \nabla \cdot v = 0 & \text{in } \Omega \end{cases}$$
(1.2)

with  $\nabla \cdot \alpha = 0$ . To describe the main theorem, we denote

$$I(x) = \int_{|y-x|<1} |v(y)|^2 dy$$

and

$$M(t) = \inf_{|x|=t} I(x).$$

Then we prove that

**Theorem 1.1.** Let  $v \in (H^1_{loc}(\Omega))^n$  be a nontrivial solution of (1.2) with an appropriate  $p \in H^1_{loc}(\Omega)$ . Assume that for  $0 \le \kappa_1 < \frac{1}{4}$ ,  $0 \le \kappa_2 < \frac{1}{2}$ ,  $0 < \delta \le \frac{1}{8}$  and  $\lambda \ge 1$ 

$$\begin{cases} |v(x)| + |\alpha(x)| + |\nabla v(x)| \le \lambda (1 + |x|^2)^{-\kappa_1 - \delta} \\ |\nabla \alpha(x)| \le \lambda (1 + |x|^2)^{-\kappa_2 - \delta}. \end{cases}$$
(1.3)

Then there exist  $\tilde{t}$  depending on  $\lambda$ , n,  $\kappa_1$ ,  $\kappa_2$ ,  $\delta$  and positive constants  $C_1$  such that

$$M(t) \ge \exp\left(-C_1 t^{\kappa} \log t\right) \quad for \quad t \ge \tilde{t}, \tag{1.4}$$

where  $\kappa = \max\{2 - 4\kappa_1, 2 - 2\kappa_2\}$  and the constant  $C_1$  depends on  $\lambda$ , n and

$$\left| \log \left( \min \{ \inf_{\tilde{t} < |x| < \tilde{t}^{(1-\delta)^{-1}}} \int_{|y-x| < 1} |v(y)|^2 dy, 1 \} \right) \right|.$$

It is interesting to compare Theorem 1.1 with the result obtained in [5] where we showed that for the standard stationary Navier-Stokes equations (i.e.,  $\alpha = 0$  in (1.2)) if v is bounded (for n = 2) or  $C^1$  bounded (for  $n \ge 3$ ) in  $\Omega$ , then

$$M(t) \ge \exp(-Ct^{2+}).$$

We can immediately deduce several consequences from Theorem 1.1. Assume that n = 3 and  $f = O(|x|^{-3})$  at infinity. Let  $u_1, u_2$  be two solutions of (1.1) satisfying  $u_1 = O(|x|^{-1})$  and  $u_2 = O(|x|^{-1})$ . It was proved by Sverak and Tsai [7] that both  $\nabla u_1$  and  $\nabla u_2$  are  $O(|x|^{-2})$ . So we can choose  $\kappa_1 = 3/16, \kappa_2 = 3/8$  (then  $\kappa = 5/4$ ), and fix  $\delta = 1/8$  in Theorem 1.1. Due to Sverak and Tsai's result, we can also relax condition (1.3). Setting  $v = u_1 - u_2$  and  $\alpha = v_1$ , we obtain from Theorem 1.1 that

**Corollary 1.2.** Let  $u_1, u_2 \in (H^1_{loc}(\Omega))^3$  be solutions of (1.1) with appropriate pressures  $p_1, p_2 \in H^1_{loc}(\Omega)$ . Assume that  $f(x) = O(|x|^{-3})$ ,  $u_1(x) = O(|x|^{-1})$ , and  $u_2 = O(|x|^{-1})$ , at infinity. Then there exist  $\tilde{t}$  and positive constant  $s_1$  such that

$$\inf_{|x|=t} \int_{|y-x|<1} |(u_1 - u_2)(y)|^2 dy \ge \exp\left(-s_1 t^{5/4} \log t\right) \quad for \quad t \ge \tilde{t}$$

where  $s_1$  depends linearly on

$$\left| \log \left( \min \{ \inf_{\tilde{t} < |x| < \tilde{t}^{8/7}} \int_{|y-x| < 1} |(u_1 - u_2)(y)|^2 dy, 1 \} \right) \right|.$$

Corollary 1.2 immediately implies the following qualitative uniqueness results.

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**Corollary 1.3.** Let  $u_1, u_2 \in (H^1_{loc}(\Omega))^3$  be solutions of (1.1) with appropriate pressures  $p_1, p_2 \in H^1_{loc}(\Omega)$ . Assume that  $f(x) = O(|x|^{-3})$ ,  $u_1(x) = O(|x|^{-1})$ , and  $u_2 = O(|x|^{-1})$ , at infinity. Then there exist R and positive constant  $s_1$  such that if

$$\int_{\Omega \cap \{|x| \ge R\}} \exp(s|x|^{5/4} \log |x|) |(u_1 - u_2)(x)|^2 dx < \infty$$

for all  $s > s_1$ , then  $u_1 \equiv u_2$  in  $\Omega$ , where  $s_1$ 's dependence is described in Corollary 1.2.

In particular, let  $u_2 = 0$  and f = 0, we have that

**Corollary 1.4.** Let n = 3, f = 0, and  $u \in (H^1_{loc}(\Omega))^3$  be a solution of (1.1) with an appropriate  $p \in H^1_{loc}(\Omega)$ . Assume that  $u(x) = O(|x|^{-1})$ . Then there exist R and positive constants  $s_1$  such that if

$$\int_{\Omega \cap \{|x| \ge R\}} \exp(s|x|^{5/4} \log |x|) |u(x)|^2 dx < \infty$$

for all  $s > s_1$ , then  $u \equiv 0$  in  $\Omega$ , where  $s_1$  depends linearly on the quantity

$$\left| \log \left( \min \{ \inf_{R < |x| < R^{\frac{8}{7}}} \int_{|y-x| < 1} |u(y)|^2 dy, 1 \} \right) \right|.$$

As in [5], we prove our result along the line of Carleman's method. Some useful techniques used in [5] are collected in the next Section. The proof of the main theorem is given in Section 3.

### 2. Reduced system and Carleman estimates

Fixing  $x_0$  with  $|x_0| = t >> 1$ , we define

$$w(x) = (at)v(atx + x_0), \ \tilde{\alpha}(x) = (at)\alpha(at + x_0), \ \text{and} \ \tilde{p}(x) = (at)^2 p(atx + x_0),$$

where  $r_1$  is the constant given in Lemma 2.1 and  $a \ge 8/r_1$  which will be determined in the proof of Theorem 1.1. Likewise, we denote

$$\Omega_t := B_{\frac{1}{a} - \frac{1}{20at^{\delta}}}(0) = \{x : |x| < \frac{1}{a} - \frac{1}{20at^{\delta}}\}.$$

From (1.2), it is easy to get that

$$\begin{cases} -\Delta w + w \cdot \nabla w + w \cdot \nabla \tilde{\alpha} + \tilde{\alpha} \cdot \nabla w + \nabla \tilde{p} = 0 & \text{in} \quad \Omega_t, \\ \nabla \cdot w = 0 & \text{in} \quad \Omega_t. \end{cases}$$
(2.1)

In view of (1.3), we have that

$$\begin{cases} \|\tilde{\alpha}\|_{L^{\infty}(\Omega_{t})} + \|w\|_{L^{\infty}(\Omega_{t})} \leq C_{0}a\lambda t^{1-2\kappa_{1}-\delta}, \\ \|\nabla w\|_{L^{\infty}(\Omega_{t})} \leq C_{0}a^{2}\lambda t^{2-2\kappa_{1}-\delta}, \\ \|\nabla \tilde{\alpha}\|_{L^{\infty}(\Omega_{t})} \leq C_{0}a^{2}\lambda t^{2-2\kappa_{2}-\frac{3}{4}\delta}, \end{cases}$$
(2.2)

where we can choose  $C_0 = (20)^{5/4}$ .

To prove Theorem1.1, we use the reduced system containing the vorticity equation derived in [5]. Let us define the vorticity q of the velocity w by

$$q = \operatorname{curl} w := \frac{1}{\sqrt{2}} (\partial_i w_j - \partial_j w_i)_{1 \le i,j \le n}.$$

The formal transpose of curl is given by

$$(\operatorname{curl}^{\top} v)_{1 \le i \le n} := \frac{1}{\sqrt{2}} \sum_{1 \le j \le n} \partial_j (v_{ij} - v_{ji}),$$

where  $v = (v_{ij})_{1 \le i,j \le n}$ . It is easy to see that

$$\Delta w = \nabla (\nabla \cdot w) - \operatorname{curl}^{\top} \operatorname{curl} w$$

(see, for example, [6] for a proof), which implies

$$\Delta w + \operatorname{curl}^{\top} q = 0 \quad \text{in} \quad \Omega_t.$$
(2.3)

Next we observe that

$$w \cdot \nabla \tilde{\alpha} + \tilde{\alpha} \cdot \nabla w = \nabla (w \cdot \tilde{\alpha}) - \sqrt{2} (\operatorname{curl} w) \tilde{\alpha} - \sqrt{2} (\operatorname{curl} \tilde{\alpha}) w$$
$$= \nabla (w \cdot \tilde{\alpha}) - \sqrt{2} q \tilde{\alpha} - \sqrt{2} (\operatorname{curl} \tilde{\alpha}) w$$

and in particular

$$w \cdot \nabla w = \nabla(\frac{1}{2}|w|^2) - \sqrt{2}(\operatorname{curl} w)w = \nabla(\frac{1}{2}|w|^2) - \sqrt{2}qw.$$

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Thus, applying curl on the first equation of (2.1), we have that

$$-\Delta q + Q(q)(w + \tilde{\alpha}) + q(\nabla w + \nabla \tilde{\alpha})^{\top} - (\nabla w + \nabla \tilde{\alpha})q^{\top} - \operatorname{div} F = 0 \text{ in } \Omega_t, \quad (2.4)$$

where

$$(Q(q)w)_{ij} = \sum_{1 \le k \le n} (\partial_j q_{ik} - \partial_i q_{jk}) w_k$$

and

$$(\operatorname{div} F)_{ij}Z = \sum_{k=1}^{n} \partial_k F_{ijk}$$

with

$$F_{ijk} = \sum_{1 \le m \le n} \left( (\operatorname{curl} \tilde{\alpha})_{jm} w_m \delta_k^i - (\operatorname{curl} \tilde{\alpha})_{im} w_m \delta_k^j \right).$$

Putting together (2.3), (2.4), and using (1.3), to prove the main theorem, it suffices to consider

$$\begin{cases} \Delta q + A(x) \cdot \nabla q + B(x)q + \operatorname{div} F = 0 & \text{in } \Omega_t, \\ \Delta w + \operatorname{curl}^\top q = 0 & \text{in } \Omega_t, \end{cases}$$
(2.5)

where A is a (3,2) tensor and B is a (2,2) tensor with

$$||A||_{L^{\infty}(\Omega_t)} \leq C_0 \lambda a t^{1-2\kappa_1-\delta}, ||B||_{L^{\infty}(\Omega_t)} \leq C_0 \lambda a^2 t^{2-2\kappa_1-\delta} + C_0 \lambda a^2 t^{2-2\kappa_2-\frac{3}{4}\delta},$$
  
and

$$|F(x)| \le C_0 \lambda a^2 t^{2-2\kappa_2 - \frac{3}{4}\delta} |w(x)|, \quad \forall \ x \in \Omega_t$$

Our proof relies on appropriate Carleman estimates. Here we need two Carleman estimates with weights  $\varphi_{\beta} = \varphi_{\beta}(x) = \exp(-\beta \tilde{\psi}(x))$ , where  $\beta > 0$ and  $\tilde{\psi}(x) = \log |x| + \log((\log |x|)^2)$ .

**Lemma 2.1.** There exist a sufficiently small number  $r_1 > 0$  depending on nand a sufficiently large number  $\beta_1 > 3$ , a positive constant C, depending on n such that for all  $v \in U_{r_1}$  and  $f = (f_1, \dots, f_n) \in (U_{r_1})^n$ ,  $\beta \ge \beta_1$ , we have that

$$\int \varphi_{\beta}^{2} (\log |x|)^{2} (\beta |x|^{4-n} |\nabla v|^{2} + \beta^{3} |x|^{2-n} |v|^{2}) dx 
\leq C \int \varphi_{\beta}^{2} (\log |x|)^{4} |x|^{2-n} [(|x|^{2} \Delta v + |x| \operatorname{div} f)^{2} + \beta^{2} ||f||^{2}] dx, \quad (2.6)$$

where  $U_{r_1} = \{ v \in C_0^{\infty}(\mathbb{R}^n \setminus \{0\}) : \operatorname{supp}(v) \subset B_{r_1} \}.$ 

Lemma 2.1 is a modified form of [4, Lemma 2.4]. For the sake of brevity, we omit the proof here. Replacing  $\beta$  of Lemma 2.1 with  $\beta + 1$  and choosing f = 0 implies

**Lemma 2.2.** There exist a sufficiently small number  $r_1 > 0$ , a sufficiently large number  $\beta_1 > 1$ , a positive constant C, such that for all  $v \in U_{r_1}$  and  $\beta \ge \beta_1$ , we have

$$\int \varphi_{\beta}^{2} (\log|x|)^{-2} |x|^{-n} (\beta|x|^{2} |\nabla v|^{2} + \beta^{3} |v|^{2}) dx \leq C \int \varphi_{\beta}^{2} |x|^{-n} (|x|^{4} |\Delta v|^{2}) dx.$$
(2.7)

In addition to Carleman estimates, we also need the following interior estimate.

**Lemma 2.3.** For any  $0 < a_1 < a_2$  such that  $B_{a_2} \subset \Omega_t$  for t > 1, let  $X = B_{a_2} \setminus \overline{B}_{a_1}$  and d(x) be the distant from  $x \in X$  to  $\mathbb{R}^n \setminus X$ . Then we have

$$\int_{X} d(x)^{2} |\nabla w|^{2} dx + \int_{X} d(x)^{4} |\nabla q|^{2} dx + \int_{X} d(x)^{2} |q|^{2} dx$$

$$\leq C \left(1 + a^{2} t^{-\frac{3\delta}{2}}\right)^{2} \int_{X} |w|^{2} dx, \qquad (2.8)$$

where the constant C depends on  $n, \lambda$ .

The proof of this lemma is similar to that given in [5].

## 3. Proof of Theorem 1.1

This section is devoted to the proof of the main theorem, Theorem 1.1. Since  $(w, p) \in (H^1(\Omega_t))^{n+1}$ , the regularity theorem implies  $w \in H^2_{loc}(\Omega_t)$ . Therefore, to use estimate (2.7), we simply cut-off w. So let  $\chi(x) \in C_0^{\infty}(\mathbb{R}^n)$ satisfy  $0 \leq \chi(x) \leq 1$  and

$$\chi(x) = \begin{cases} 0, & |x| \le \frac{1}{8at}, \\ 1, & \frac{1}{4at} < |x| < \frac{1}{a} - \frac{3}{20at^{\delta}}, \\ 0, & |x| \ge \frac{1}{a} - \frac{2}{20at^{\delta}}. \end{cases}$$

It is easy to see that for any multiindex  $\alpha$ 

$$\begin{cases} |D^{\alpha}\chi| = O((at)^{|\alpha|}) & \text{if } \frac{1}{8at} \le |x| \le \frac{1}{4at}, \\ |D^{\alpha}\chi| = O((at^{\delta})^{|\alpha|}) & \text{if } \frac{1}{a} - \frac{3}{20at^{\delta}} \le |x| \le \frac{1}{a} - \frac{2}{20at^{\delta}}. \end{cases}$$
(3.1)

To apply Carleman estimates above, it suffices to take  $1/a \leq r_1$ . Now applying (2.7) to  $\chi w$  gives

$$\int (\log|x|)^{-2} \varphi_{\beta}^{2} |x|^{-n} (\beta|x|^{2} |\nabla(\chi w)|^{2} + \beta^{3} |\chi w|^{2}) dx$$

$$\leq C \int \varphi_{\beta}^{2} |x|^{-n} |x|^{4} |\Delta(\chi w)|^{2} dx. \qquad (3.2)$$

Here and after, C and  $\tilde{C}$  denote general constants whose value may vary from line to line. The dependence of C and  $\tilde{C}$  will be specified whenever necessary. Next applying (2.6) to  $v = \chi q$  and  $f = |x|\chi F$  yields that

$$\int \varphi_{\beta}^{2} (\log |x|)^{2} (|x|^{4-n}\beta |\nabla(\chi q)|^{2} + |x|^{2-n}\beta^{3} |\chi q|^{2}) dx$$
  
$$\leq C \int \varphi_{\beta}^{2} (\log |x|)^{4} |x|^{2-n} [(|x|^{2}\Delta(\chi q) + |x|\operatorname{div}(|x|\chi F))^{2} + \beta^{2} ||x|\chi F||^{2}] dx. (3.3)$$

Combining  $\beta \times (3.2)$  and (3.3), we obtain that

$$\int_{W} (\log |x|)^{-2} \varphi_{\beta}^{2} |x|^{-n} (\beta^{2} |x|^{2} |\nabla w|^{2} + \beta^{4} |w|^{2}) dx 
+ \int_{W} (\log |x|)^{2} \varphi_{\beta}^{2} |x|^{-n} (\beta |x|^{4} |\nabla q|^{2} + |x|^{2} \beta^{3} |q|^{2}) dx 
\leq C\beta \int \varphi_{\beta}^{2} |x|^{-n} |x|^{4} |\Delta(\chi w)|^{2} dx 
+ C \int \varphi_{\beta}^{2} (\log |x|)^{4} |x|^{2-n} [(|x|^{2} \Delta(\chi q) + |x| \operatorname{div}(|x|\chi F))^{2} 
+ \beta^{2} |||x| \chi F ||^{2}] dx,$$
(3.4)

where W denotes the domain  $\{x : \frac{1}{4at} < |x| < \frac{1}{a} - \frac{3}{20at^{\delta}}\}$ . To simplify the notations, we denote  $Y = \{x : \frac{1}{8at} \le |x| \le \frac{1}{4at}\}$  and  $Z = \{x : \frac{1}{a} - \frac{3}{20at^{\delta}} \le |x| \le \frac{1}{a} - \frac{2}{20at^{\delta}}\}$ . By (2.4) and estimates (3.1), we deduce from (3.4) that

$$\int_{W} (\log |x|)^{-2} \varphi_{\beta}^{2} |x|^{-n} (\beta^{2} |x|^{2} |\nabla w|^{2} + \beta^{4} |w|^{2}) dx$$

$$+ \int_{W} (\log |x|)^{2} \varphi_{\beta}^{2} |x|^{-n} (\beta |x|^{4} |\nabla q|^{2} + |x|^{2} \beta^{3} |q|^{2}) dx$$

$$\leq C\beta \int_{W} \varphi_{\beta}^{2} |x|^{-n} |x|^{4} |\nabla q|^{2} dx$$

$$+ Ca^{2} t^{2-4\kappa_{1}-2\delta} \int_{W} (\log |x|)^{4} \varphi_{\beta}^{2} |x|^{-n} |x|^{6} |\nabla q|^{2} dx$$

$$+ Ca^{4} t^{4-4\kappa_{1}-2\delta} \int_{W} (\log |x|)^{4} \varphi_{\beta}^{2} |x|^{-n} |x|^{6} |q|^{2} dx$$

$$+ C\beta^{2} a^{4} t^{4-4\kappa_{2}-\frac{3}{4}\delta} \int_{W} (\log |x|)^{4} \varphi_{\beta}^{2} |x|^{-n} |x|^{4} |w|^{2} dx$$

$$+ C(at)^{4} \beta \int_{Y \cup Z} \varphi_{\beta}^{2} |x|^{-n} |\tilde{U}|^{2} dx$$

$$+ C(at)^{4} \beta^{2} \int_{Y \cup Z} (\log |x|)^{4} \varphi_{\beta}^{2} |x|^{2-n} |\tilde{U}|^{2} dx,$$

$$(3.5)$$

where  $|\tilde{U}(x)|^2 = |x|^4 |\nabla q|^2 + |x|^2 |q|^2 + |x|^2 |\nabla w|^2 + |w|^2$  and C depends on n,  $\lambda$ .

Now we can choose  $a > a_0 \ge 8/r_1$  such that  $(\log |x|)^2 \ge 2C$  for all  $x \in W$ . Then the first term on the right hand side of (3.5) can be absorbed by the left hand side of (3.5). Now, let  $\beta \ge \beta_2 = t^{\kappa}$  and choose  $t \ge t_0$  with  $t_0$  depending on  $a, \lambda, \delta$  such that the second term to the fourth term on the right hand side of (3.5) can be removed. With the choices described above, we obtain from (3.5) that

$$\begin{split} \beta^{4}(b_{1})^{-n}(\log b_{1})^{-2}\varphi_{\beta}^{2}(b_{1}) \int_{\frac{1}{at} < |x| < b_{1}} |w|^{2} dx \\ &\leq \beta^{4} \int_{W} (\log |x|)^{-2}\varphi_{\beta}^{2} |x|^{-n} |w|^{2} dx \\ &\leq C\beta(at)^{4} \int_{Y \cup Z} (\log |x|)^{4}\varphi_{\beta}^{2} |x|^{-n} |\tilde{U}|^{2} dx \\ &\leq C\beta^{2}(at)^{4} (\log b_{2})^{4} b_{2}^{-n} \varphi_{\beta}^{2}(b_{2}) \int_{Y} |\tilde{U}|^{2} dx \\ &+ C\beta^{2}(at)^{4} (\log b_{3})^{4} b_{3}^{-n} \varphi_{\beta}^{2}(b_{3}) \int_{Z} |\tilde{U}|^{2} dx, \end{split}$$
(3.6)

where  $b_1 = \frac{1}{a} - \frac{8}{20at^{\delta}}$ ,  $b_2 = \frac{1}{8at}$  and  $b_3 = \frac{1}{a} - \frac{3}{20at^{\delta}}$ .

Using (2.8), we can control  $|\tilde{U}|^2$  terms on the right hand side of (3.6).

Indeed, let  $X = Y_1 := \{x : \frac{1}{16at} \le |x| \le \frac{1}{2at}\}$ , then we can see that

$$d(x) \ge C|x|$$
 for all  $x \in Y$ ,

where C an absolute constant. Therefore, (2.8) implies

$$\int_{Y} \left( |x|^{2} |\nabla w|^{2} + |x|^{4} |\nabla q|^{2} + |x|^{2} |q|^{2} \right) dx 
\leq C \int_{Y_{1}} \left( d(x)^{2} |\nabla w|^{2} + d(x)^{4} |\nabla q|^{2} + d(x)^{2} |q|^{2} \right) dx 
\leq C \left( 1 + a^{2} t^{-\frac{3\delta}{2}} \right)^{2} \int_{Y_{1}} |w|^{2} dx 
\leq C a^{4} \int_{Y_{1}} |w|^{2} dx.$$
(3.7)

Here C depends on  $n, \lambda$ . On the other hand, let  $X = Z_1 := \{x : \frac{1}{2a} \le |x| \le \frac{1}{a} - \frac{1}{20at^{\delta}}\}$ , then

$$d(x) \ge Ct^{-\delta}|x|$$
 for all  $x \in Z$ ,

where C another absolute constant. Thus, it follows from (2.8) that

$$\int_{Z} \left( |x|^{2} |\nabla w|^{2} + |x|^{4} |\nabla q|^{2} + |x|^{2} |q|^{2} \right) dx 
\leq Ct^{4\delta} \int_{Z_{1}} \left( d(x)^{2} |\nabla w|^{2} + d(x)^{4} |\nabla q|^{2} dx + d(x)^{2} |q|^{2} \right) dx 
\leq Ct^{4\delta} \left( 1 + a^{2} t^{-\frac{3\delta}{2}} \right)^{2} \int_{Z_{1}} |w|^{2} dx 
\leq C(at)^{4} \int_{Z_{1}} |w|^{2} dx.$$
(3.8)

Combining (3.6), (3.7), and (3.8) leads to

$$b_{1}^{-2\beta-n} (\log b_{1})^{-4\beta-2} \int_{\frac{1}{2at} < |x| < b_{1}} |w|^{2} dx$$

$$\leq Ca^{8} t^{4} (\log b_{2})^{4} b_{2}^{-n} \varphi_{\beta}^{2} (b_{2}) \int_{Y_{1}} |w|^{2} dx$$

$$+ C(at)^{8} (\log b_{3})^{4} b_{3}^{-n} \varphi_{\beta}^{2} (b_{3}) \int_{Z_{1}} |w|^{2} dx.$$
(3.9)

Notice that (3.9) holds for all  $\beta \geq \beta_2$ .

Changing  $2\beta + n$  to  $\beta$ , (3.9) becomes

$$b_{1}^{-\beta} (\log b_{1})^{-2\beta+2n-2} \int_{\frac{1}{2at} < |x| < b_{1}} |w|^{2} dx$$

$$\leq Ca^{8} t^{4} b_{2}^{-\beta} (\log b_{2})^{-2\beta+2n+4} \int_{Y_{1}} |w|^{2} dx$$

$$+ C(at)^{8} b_{3}^{-\beta} (\log b_{3})^{-2\beta+2n+4} \int_{Z_{1}} |w|^{2} dx. \qquad (3.10)$$

Dividing  $b_1^{-\beta}(\log b_1)^{-2\beta+2n-2}$  on the both sides of (3.10) and noting  $\beta \ge n+2 > n-1$ , i.e.,  $2\beta - 2n + 2 > 0$ , we have for  $t \ge t_1 \ge t_0$  that

$$\int_{|x+\frac{b_{4}x_{0}}{t}|<\frac{1}{at}} |w(x)|^{2} dx 
\leq \int_{\frac{1}{2at}<|x|
(3.11)$$

where  $b_4 = \frac{1}{a} - \frac{1}{at^{\delta}}$  and  $b_5 = \frac{1}{a} - \frac{6}{20at^{\delta}}$ . In deriving the third inequality above, we use the fact that

$$0 \le (\frac{b_5}{b_3})(\frac{\log b_1}{\log b_3})^2 = 1 - \frac{1}{2t^{\delta}\log a} - \frac{3}{20t^{\delta}} + O(t^{-2\delta}) \le 1$$

for all  $t \ge t_2 \ge t_1$  and  $a > a_1 = \max\{1, a_0\}$ , where  $t_2$  depends on  $t_1$ ,  $\delta$ , and a. From now on we fix a, which depends only on n and  $r_1$ . Recall that  $r_1$  is a function of n. Therefore,  $t_2$  depends on n,  $\lambda$ , and  $\delta$ . Having fixed constant a,  $|\log b_3|$  can be bounded by a positive constant. Thus, (3.11) is reduced to

$$\int_{|x+\frac{b_4x_0}{t}|<\frac{1}{at}} |w(x)|^2 dx \leq Ct^4 (\log t)^6 (8t)^\beta \int_{|x|<\frac{1}{at}} |w(x)|^2 dx$$

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$$+Ct^8(b_1/b_5)^\beta \int_{Z_1} |w(x)|^2 dx, \qquad (3.12)$$

where C depends on n and  $\lambda$ .

From (3.12), (2.2), the definition of w(x), the change of variables  $y = atx + x_0$ , and  $x_0 = ty_0$ , we have that

$$I(t^{1-\delta}y_0) \leq Ct^4 (\log t)^6 (8t)^{\beta} \int_{|y-x_0|<1} |u(y)|^2 dy + Ct^{8-\frac{3\delta}{2}} \left(\frac{t^{\delta}}{t^{\delta} + \frac{1}{10}}\right)^{\beta}$$
  
$$\leq C(8t)^{\beta+10} I(ty_0) + Ct^8 \left(\frac{t^{\delta}}{t^{\delta} + \frac{1}{10}}\right)^{\beta}$$
  
$$\leq C(8t)^{2\beta} I(ty_0) + Ct^8 \left(\frac{t^{\delta}}{t^{\delta} + \frac{1}{10}}\right)^{\beta}$$
(3.13)

provided  $\beta \geq \beta_2$ . For simplicity, by denoting

$$A(t) = 2\log 8t, \quad B(t) = \log(\frac{t^{\delta} + \frac{1}{10}}{t^{\delta}}),$$

(3.13) becomes

$$I(t^{1-\delta}y_0) \le C \Big\{ \exp(\beta A(t))I(ty_0) + t^8 \exp(-\beta B(t)) \Big\}.$$
 (3.14)

Now, we consider two cases. If

$$\exp(\beta_2 A(t))I(ty_0) \ge t^8 \exp(-\beta_2 B(t)),$$

then we have

$$I(x_0) = I(ty_0) \ge t^8 \exp(-\beta_2(A(t) + B(t))) = t^8(8t)^{-2\beta_2} \left(\frac{t^{\delta} + \frac{1}{10}}{t^{\delta}}\right)^{-\beta_2},$$

that is

$$I(ty_0) \ge t^{-2\beta_2 + 8} = t^{-2t^{\kappa} + 8} \ge \exp(-2t^{\kappa}\log t)$$
(3.15)

for any fixed  $t \ge t_2$ . Note that we have used the relation  $\beta_2 = t^{\kappa}$  in (3.15).

On the other hand, if

$$\exp(\beta_2 A(t))I(ty_0) < t^8 \exp(-\beta_2 B(t)),$$

then we can pick a  $\tilde{\beta} > \beta_2$  such that

$$\exp(\tilde{\beta}A(t))I(ty_0) = t^8 \exp(-\tilde{\beta}B(t)).$$
(3.16)

Solving  $\tilde{\beta}$  from (3.16) and using (3.14), we have that

$$I(t^{1-\delta}y_0) \leq C \exp(\tilde{\beta}A(t))I(ty_0) = C (I(ty_0))^{\tau} (t^8)^{1-\tau} \leq C t^8 (I(ty_0))^{\tau}, \qquad (3.17)$$

where  $\tau = \frac{B(t)}{A(t) + B(t)}$ .

It is time to prove Theorem 1.1. Let  $|x_0| = t$  for  $t \ge t_2^{\frac{1}{1-\delta}}$  and  $y_0 = \frac{x_0}{t}$ , then we can write

$$t = \mu^{\left((1-\delta)^{-s}\right)} \tag{3.18}$$

for some positive integer s and  $t_2 \leq \mu < t_2^{\frac{1}{1-\delta}} \leq t_2^2$ . For simplicity, we define  $d_j = \mu^{\left((1-\delta)^{-j}\right)}$  and  $\tau_j = \frac{B(d_j)}{A(d_j)+B(d_j)}$  for  $j = 1, 2 \cdots s$ . Define

$$J = \{1 \le j \le s : \exp(d_j^{\kappa} A(d_j)) I(d^j y_0) \ge d_j^8 \exp(-d_j^{\kappa} B(d_j))\}.$$

Now, we divide it into two cases. If  $J = \emptyset$ , we only need to consider (3.17). Using (3.17) iteratively starting from  $t = d_1$ , we have that

$$I(\mu y_0) \leq C(d_1^8) (I(d_1 y_0))^{\tau_1} \\ \leq C^s (d_1 d_2 \cdots d_s)^8 (I(x_0))^{\tau_1 \tau_2 \cdots \tau_s}.$$
(3.19)

By (3.18) and (3.19), we obtain that

$$I(\mu y_0) \leq C^{(\log \log t/|\log(1-\delta)|)} t^{8/\delta} (I(x_0))^{\tau_1 \tau_2 \cdots \tau_s} \leq t^{\tilde{C}_0/\delta} (I(x_0))^{\tau_1 \tau_2 \cdots \tau_s},$$
(3.20)

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where  $\tilde{C}_0$  depends on  $\lambda$ , n. It is easily to see that

$$\frac{1}{\tau_j} = \frac{2\log(8d_j) + \log(1 + 0.1d_j^{-\delta})}{\log(1 + 0.1d_j^{-\delta})} \le \frac{4\log(8d_j)}{\log(1 + 0.1d_j^{-\delta})} \le 160d_j^{\delta}\log(d_j),$$

and thus

$$\frac{1}{\tau_1 \tau_2 \cdots \tau_s} \leq (160 \log \mu \log t)^s (d_1 \cdots d_s)^\delta \\ \leq t \omega(t), \tag{3.21}$$

where  $\omega(t) = (\log t)^{4\log(\log t)}$ . Raising both sides of (3.20) to the power  $\frac{1}{\tau_1 \tau_2 \cdots \tau_s}$  and using (3.21), we obtain that

$$(\min\{I(\mu y_0), 1\})^{t\omega(t)} \leq I(\mu y_0)^{\frac{1}{\tau_1 \tau_2 \cdots \tau_s}} \leq e^{(\tilde{C}_0/\delta)t\omega(t)} (I(x_0)).$$
(3.22)

Next, if  $J \neq \emptyset$ , let *l* be the largest integer in *J*. Then from (3.15) we have

$$I(d_l y_0) \ge d_l^{-2d_l^{\kappa}+8}.$$
 (3.23)

Iterating (3.17) starting from  $t = d_{l+1}$  yields

$$I(d_{l}y_{0}) \leq C^{s-l}(d_{l+1}\cdots d_{s})^{8} (I(x_{0}))^{\tau_{l+1}\cdots\tau_{s}}$$
  
$$\leq C^{(\log\log t/|\log(1-\delta)|)}(t/d_{l})^{8/\delta} (I(x_{0}))^{\tau_{l+1}\cdots\tau_{s}}$$
  
$$\leq t^{\tilde{C}_{0}/\delta} (I(x_{0}))^{\tau_{l+1}\cdots\tau_{s}}.$$
(3.24)

It is enough to assume  $I(d_l y_0) < 1$ . Repeating the computations in (3.21), we can see that

$$\frac{1}{\tau_{l+1}\cdots\tau_s} \le (t/d_l)\omega(t). \tag{3.25}$$

Hence, combining (3.23), (3.24) and using (3.25), we get that

$$t^{-\hat{C}_{3}t^{\kappa}\log(t)} \le e^{(\hat{C}_{0}/\delta)t\omega(t)} \left(I(x_{0})\right), \qquad (3.26)$$

where  $\tilde{C}_3$  is an absolute constant. The proof is complete in view of (3.15), (3.22) and (3.26).

## Acknowledgments

The authors were supported in part by the Ministry of Science and Technology, Taiwan.

# References

- R. Finn, Stationary solutions of the Navier-Stokes equations, Proc. Symp. Appl. Math. Amer. Math. Soc., 17 (1965), 121-153.
- R. Finn, On steady-state solutions for the Navier-Stokes partial differential equations, Arch. Rational Mech. Anal., 3 (1959), 381-396.
- L. Hörmander, The analysis of linear partial differential operators, Vol. 3, Springer-Verlag, Berlin/New York, 1985.
- C.L. Lin, G. Nakamura and J.N. Wang, Optimal three-ball inequalities and quantitative uniqueness for the Lamé system with Lipschitz coefficients, *Duke Math Journal*, 155 (2010), No. 1, 189-204.
- C. L. Lin, G. Uhlmann and J. N. Wang, Asymptotic behavior of solutions of the stationary Navier-Stokes equations in an exterior domain, *Indiana University Mathematics Journal*, **60** (2011), No. 6, 2093-2106.
- 6. M. Mitrea and S. Monniaux, Maximal regularity for the Lamé system in certain classes of non-smooth domains, J. Evol. Equ., DOI 10.1007/s00028-010-0071-1.
- V. Sverak and T. P. Tsai, On the spatial decay of 3-D steady-state Navier-Stokes flows, Comm in PDE, 25 (2000), 2107-2117.